On smooth approximations of rough vector fields and the selection of flows

G. Ciampa, G. Crippa and S. Spirito

Departement Mathematik und Informatik Fachbereich Mathematik Universität Basel CH-4051 Basel Preprint No. 2019-04 Februar 2019

www.math.unibas.ch

ON SMOOTH APPROXIMATIONS OF ROUGH VECTOR FIELDS AND THE SELECTION OF FLOWS

GENNARO CIAMPA, GIANLUCA CRIPPA, AND STEFANO SPIRITO

ABSTRACT. In this work we deal with the selection problem of flows of an irregular vector field. We first summarize an example from [4] of a vector field b and a smooth approximation b_{ε} for which the sequence X^{ε} of flows of b_{ε} has subsequences converging to different flows of the limit vector field b. Furthermore, we give some heuristic ideas on the selection of a subclass of flows in our specific case.

1. Introduction and notations

Consider the system of ordinary differential equations

$$\begin{cases} \frac{\mathrm{d}}{\mathrm{d}t} X(t,x) = b(t,X(t,x)), \\ X(0,x) = x, \end{cases}$$
 (1.1)

where $(t,x) \in (0,T) \times \mathbb{R}^d$ are the independent variables, with $T < \infty$, $b:(0,T) \times \mathbb{R}^d \to \mathbb{R}^d$ is a given vector field and $X:(0,T) \times \mathbb{R}^d \to \mathbb{R}^d$ is the unknown. A solution X of (1.1) is called flow of b. The well posedness of (1.1) is a well known result when the vector field b is globally Lipschitz in space uniformly in time. The system (1.1) is strictly connected to the Cauchy problem for the linear transport equation

$$\begin{cases} \partial_t u + b \cdot \nabla u = 0, \\ u|_{t=0} = u_0, \end{cases}$$
 (1.2)

since in a smooth setting, the unique solution of (1.2) is given by the formula $u(t,x) = u_0((X(t,\cdot))^{-1}(x))$, where X is the unique flow of b.

Besides the theoretical interest, due to applications to several equations from mathematical physics the setting of smooth vector fields is too restrictive and a theory under assumptions of lower regularity was developed in the last years. Exploiting the connection between (1.2) and (1.1), DiPerna and Lions in [8] proved the well posedness of (1.2) under hypotesis of Sobolev regularity for the vector field and bounded divergence. As a consequence of their result, they proved well posedness of (1.1) under the same hypotesis. Similarly, Ambrosio in [1] improved the result of [8] to the case of BV regularity and bounded divergence for b. On the other hand, a well posedness theory based only on a prori estimates of the flow was developed in [5] for $W^{1,p}$ vector field with p > 1 and in [3],[6] for the case p = 1 and vector fields whith gradient given by a singular integral of a L^1 function. This latter is a class of interest in the context of 2D Euler equations. More recently Nguyen in [9] improved the result to vector fields which can be represented as singular integral of a function in BV.

Various counterexamples show that weak differentiability assumptions on the vector field are in general necessary in order to obtain well posedness, see for instance [7], [8]. For a general survey on this topic, we refer to [2]. The aim of this note is to discuss the selection problem for solutions of (1.1) in a low regularity setting. To better explain what we mean by selection, let us first recall some preliminary notations and definitions.

We denote by \mathcal{L}^d the Lebesgue measure on \mathbb{R}^d . If $f: \mathbb{R}^d \to \mathbb{R}$ is a Borel map we denote by $f \# \mathcal{L}^d$ the *push forward*, that is, the measure defined by the following relation

$$f # \mathcal{L}^d(E) = \mathcal{L}^d(f^{-1}(E))$$
 for every Borel set $E \subset \mathbb{R}^d$.

The definition of flow of a vector field b, when b is not smooth, is the following:

Definition 1.1. Let $b \in L^1((0,T); L^1_{loc}(\mathbb{R}^d; \mathbb{R}^d))$ be given. We say that $X: (0,T) \times \mathbb{R}^d \to \mathbb{R}^d$ is a regular Lagrangian flow associated to b if

(1) for a.e. $x \in \mathbb{R}^d$ the map $t \to X(t,x)$ is an absolutely continuous integral solution of the ordinary differential equation

$$\begin{cases} \frac{\mathrm{d}}{\mathrm{d}t} X(t,x) = b(t,X(t,x)), \\ X(0,x) = x, \end{cases}$$
 (1.3)

(2) there exists a constant L indipendent of t such that

$$X(t,\cdot)\#\mathcal{L}^d \le L\mathcal{L}^d. \tag{1.4}$$

If the vector field is divergence-free, L can be taken to be 1 and (1.4) is an equality. This means that the flow X is measure preserving. Condition (1.4) is a first selection: we only consider among all solutions of (1.1) those that do not compress too much the Lebesgue measure. This selection is necessary in the theory since it is not known if there is uniqueness in the class of flows that can compress the Lebesgue measure, even under assumptions of weak differentiability and zero divergence for the vector field.

The paper is divided as follows. In Section 2 we give a precise statement for the selection problem and we give an example of a vector field and of an approximation for which the selection is not true. In Section 3 we characterize a class of flows through measure preserving maps of the unit circle. Finally in Section 4 we introduce a new question about the selection of a subset of the set of all flows and we give some ideas and heuristics about what we can expect.

2. The problem of selection

Let us consider a weakly differentiable vector field b which falls into the class of well-posedness like those discussed in the introduction. To prove the existence of solutions of (1.1), the natural approach is to rely on a compactness argument for an approximating sequence X^{ε} . This latter is usually constructed as the (unique) flow of a smooth approximation b_{ε} of b, see [2]. Consider, instead, a vector field b that has more than one regular Lagrangian flow and let b_{ε} be a smooth approximation of b. Consider the solution X^{ε} of the ODE relative to b_{ε} and assume that X^{ε} converges to a

regular Lagrangian flow X of b. We wonder if for every approximation b_{ε} the corresponding flows X^{ε} can converge to only one regular Lagrangian flow: if this were true, this procedure could be considered as a selection principle for the flows of an irregular vector field. We can summarize the previous discussion in the following: question

(Q1) Does the approximation procedure obtained by smoothing the vector field select a unique solution of (1.1)?

In [4] we give a negative answer to the previous question showing a counterexample. Precisely, we consider this vector field, which is a 3D analogous of an example of DiPerna and Lions [8]:

$$b(x,y,z) = \begin{cases} \left(-\operatorname{sgn}(z)\frac{x}{|z|^2}, -\operatorname{sgn}(z)\frac{y}{|z|^2}, -\frac{2}{|z|}\right) & \text{if } x \in P, \\ (0,0,0) & \text{otherwise,} \end{cases}$$
(2.1)

where $P \subset \mathbb{R}^3$ denotes the set

$$P = P^{+} \cup P^{-} = \{(x, y, z) \in \mathbb{R}^{3} : x^{2} + y^{2} \le z\} \cup \{(x, y, z) \in \mathbb{R}^{3} : x^{2} + y^{2} \le -z\},\$$

the union of two symmetric paraboloids.

The vector field b is divergence-free and it is out of the class of uniqueness of solutions of (1.1). In particular, observe that b is not in any Sobolev space $W^{1,p}(\mathbb{R}^3)$ or in $BV(\mathbb{R}^3)$.

We want to define two different regular Lagrangian flows \bar{X}, \tilde{X} of b and, since we are considering flows defined almost everywhere, we need to define \bar{X}, \tilde{X} only on $\mathbb{R}^3 \setminus \{0\}$. We start for $\mathbf{x} \in \mathbb{R}^3 \setminus P$: in this region the vector field is identically 0 so that we define \bar{X}, \tilde{X} simply as

$$\bar{X}(t, \mathbf{x}) = \mathbf{x} = \tilde{X}(t, \mathbf{x}) \qquad \forall t \ge 0.$$

If $\mathbf{x} = (x, y, z) \in P^-$ we define \bar{X}, \tilde{X} as

$$\begin{cases} \bar{X}_{1}(t,x,z) = \tilde{X}_{1}(t,x,z) = \frac{x}{\sqrt{-z}} \sqrt[4]{z^{2} + 4t} \\ \bar{X}_{2}(t,y,z) = \tilde{X}_{2}(t,y,z) = \frac{y}{\sqrt{-z}} \sqrt[4]{z^{2} + 4t} \end{cases} \quad \forall t \geq 0.$$

$$(2.2)$$

$$\bar{X}_{3}(t,z) = \tilde{X}_{3}(t,z) = -\sqrt{z^{2} + 4t}$$

Finally, when $\mathbf{x} = (x, y, z) \in P^+$ define the flows as

$$\begin{cases}
\bar{X}_{1}(t, x, z) = \tilde{X}_{1}(t, x, z) = \frac{x}{\sqrt{z}} \sqrt[4]{z^{2} - 4t} \\
\bar{X}_{2}(t, y, z) = \tilde{X}_{2}(t, y, z) = \frac{y}{\sqrt{z}} \sqrt[4]{z^{2} - 4t} & \text{for } t \in \left[0, \frac{z^{2}}{4}\right]. \\
\bar{X}_{3}(t, z) = \tilde{X}_{3}(t, z) = \sqrt{z^{2} - 4t}
\end{cases}$$
(2.3)

At time $t = \frac{z^2}{4}$ the trajectories reach the origin and then one possible way to extend them for later times is

extend them for later times is
$$\begin{cases}
\bar{X}_{1}(t, x, z) = \frac{x}{\sqrt{z}} \sqrt[4]{4t - z^{2}} \cos \Theta - \frac{y}{\sqrt{z}} \sqrt[4]{4t - z^{2}} \sin \Theta \\
\bar{X}_{2}(t, y, z) = \frac{x}{\sqrt{z}} \sqrt[4]{4t - z^{2}} \sin \Theta + \frac{y}{\sqrt{z}} \sqrt[4]{4t - z^{2}} \cos \Theta & t \geq \frac{z^{2}}{4}, \\
\bar{X}_{3}(t, z) = -\sqrt{4t - z^{2}}
\end{cases}$$
(2.4)

while

$$\begin{cases} \tilde{X}_1(t,x,z) = \frac{x}{\sqrt{z}} \sqrt[4]{4t - z^2} \cos \Phi - \frac{y}{\sqrt{z}} \sqrt[4]{4t - z^2} \sin \Phi \\ \tilde{X}_2(t,y,z) = \frac{x}{\sqrt{z}} \sqrt[4]{4t - z^2} \sin \Phi + \frac{y}{\sqrt{z}} \sqrt[4]{4t - z^2} \cos \Phi \end{cases} \quad t \ge \frac{z^2}{4}, \quad (2.5)$$

where $\Theta, \Phi \in (0, 2\pi]$ and $\Theta \neq \Phi$. An easy computation shows that \bar{X}, \tilde{X} are two different regular Lagrangian flows of b. We call those kind of solutions respectively X^{Θ}, X^{Φ} , where Θ and Φ represent a rotation in the xy plane. Heuristically, we can define this kind of flows as a consequence of the fact that the trajectories once they reach the origin can come out arbitrarily. In [4] one of our main results is the following:

Theorem 2.1. There exists a sequence of vector fields $b_n \in C^{\infty}(\mathbb{R}^3)$ such that:

- b_n is divergence-free;
- b_n → b in L¹_{loc}(ℝ³);
 the sequence Xⁿ of regular Lagrangian flows of b_n has two different subsequences converging in $L^{\infty}((0,T);L^1_{loc}(\mathbb{R}^3)))$ to two different regular Lagrangian flows of b.

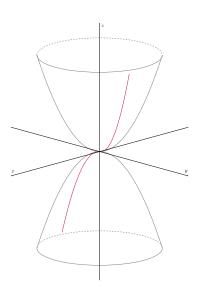


FIGURE 1. An example of solution X^{Θ}

In the proof of Theorem 2.1, given $\Theta, \Phi \in (0, 2\pi]$ with $\Theta \neq \Phi$, we construct an explicit approximation which has two different subsequences converging respectively to X^{Θ} and X^{Φ} . The strategy of the approximation is based on smoothing b nearby the origin and forcing the trajectories to rotate very fast along a given helix. We basically modify b in a small region with contains the singularity and leave the rest unchanged.

The theorem answers question (Q1) in the negative. However with our

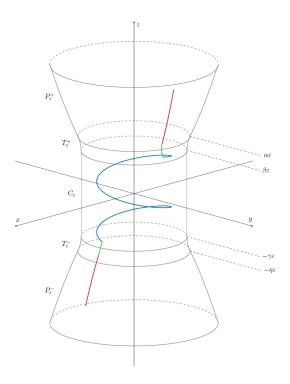


FIGURE 2. The figure represents an approximated trajectory in the construction of the proof of Theorem 2.1

approach we are able to obtain only solutions of the form X^{Θ} . Indeed, note that another possible way to define a flow for $\mathbf{x} \in P^+$ is the following:

her possible way to define a flow for
$$\mathbf{x} \in P^+$$
 is the following:
$$\begin{cases} X_1(t,r,\theta,z) = \frac{r}{\sqrt{z}} \sqrt[4]{z^2 - 4t} \cos \theta \\ X_2(t,r,\theta,z) = \frac{r}{\sqrt{z}} \sqrt[4]{z^2 - 4t} \sin \theta & \text{for } t \in \left[0, \frac{z^2}{4}\right], \\ X_3(t,z) = \sqrt{z^2 - 4t} \end{cases}$$
(2.6)

and

$$\begin{cases} X_1(t, r, \theta, z) = \frac{r}{\sqrt{z}} \sqrt[4]{4t - z^2} \cos \psi(\theta) \\ X_2(t, r, \theta, z) = \frac{r}{\sqrt{z}} \sqrt[4]{4t - z^2} \sin \psi(\theta) & \text{for } t \ge \frac{z^2}{4}, \\ X_3(t, z) = -\sqrt{4t - z^2} \end{cases}$$
 (2.7)

where the map $\psi:[0,2\pi] \to [0,2\pi]$ is arbitrary and (r,θ,z) denote the cylindrical coordinates in \mathbb{R}^3 . It is easy to check that the map in (2.6),(2.7) is a solution of the ODE relative to b; we call X_{ψ} such a map. It will turn out to be useful the flow on $P \setminus \{0\}$ and not only in \mathring{P} although we deal with functions defined almost everywhere with respect to the 3D Lebesgue measure. The reason for that lies in the fact that for our purpose we will compute X on $\partial P \setminus \{0\}$; this would not make sense without a suitable definition of X on the boundary of P. Such definition is made accordingly to the everywhere definition of b. In the next section we will discuss the conditions that the

map ψ has to satisfied in order for X_{ψ} to be a regular Lagrangian flow of b.

3. Regular Lagrangian flows and measure preserving map on the circle

In this section we prove that solutions of the form X_{ψ} are regular Lagrangian flows of b when ψ is a measure preserving map. Before doing this, note that the map X_{ψ} associated to $\psi(\theta) = \alpha$, where $\alpha \in (0, 2\pi]$ is fixed, is a solution of the ODE but it does not preserve the 3D Lebesgue measure and then it is not a regular Lagrangian flow.

Now we recall the definition of a mesure preserving map on the unit circle.

Definition 3.1. Let $\psi : \mathbb{S}^1 \to \mathbb{S}^1$ be a measurable map, where $\mathbb{S}^1 = \mathbb{R}/2\pi\mathbb{Z}$ is the unit circle with the 1D Lebesgue measure. The map ψ is called measure preserving if

$$\psi \# \mathscr{L}^1 = \mathscr{L}^1.$$

We identify $\mathbb{S}^1 \sim [0, 2\pi]$ and we define the set \mathcal{M} as

$$\mathcal{M} := \{ \psi : [0, 2\pi] \to [0, 2\pi] : \psi \text{ satisfies Definition 3.1} \}.$$

Moreover, define the maps

$$I_{\pm}: \theta \in [0, 2\pi] \to (\cos \theta, \sin \theta, \pm 1) \in \mathbb{R}^3.$$

Proposition 3.2. Given a regular Lagrangian flow X there exists $\psi \in \mathcal{M}$ such that $X = X_{\psi}$. Viceversa given $\psi \in \mathcal{M}$ there exists a unique regular Lagrangian flow X such that $X = X_{\psi}$.

Proof. Consider a regular Lagrangian flow $X(t, \mathbf{x})$ and define

$$\psi(\theta) = I_{-}^{-1} \left(X \left(\frac{1}{2}, I_{+}(\theta) \right) \right) \qquad \theta \in [0, 2\pi].$$

We need to show that such a map preserves the 1D Lebesgue measure: consider a Borel set $E \subseteq [0, 2\pi]$ and define **E** as the set

$$\mathbf{E} = \{ (\rho, \theta, z) : \theta \in E, \rho \in [0, \sqrt{z}], z \in [-1, 0] \}.$$

A straightforward computation shows that

$$X^{-1}\left(\frac{1}{2}, \cdot\right)(\mathbf{E}) = \{(\rho, \theta, z) : \theta \in \psi^{-1}(E), \rho \in [0, \sqrt{z}], z \in [1, \sqrt{2}]\}$$

and

$$\mathscr{L}^{1}(\psi^{-1}(E)) = 4\mathscr{L}^{3}\left(X^{-1}\left(\frac{1}{2},\cdot\right)(\mathbf{E})\right) = 4\mathscr{L}^{3}(\mathbf{E}) = \mathscr{L}^{1}(E), \quad (3.1)$$

hence ψ is measure preserving.

We now prove the other implication. Consider a measure preserving map ψ , a point $\mathbf{x} \in \mathbb{S}^1 \times \{1\}$ and solve the system

$$\begin{cases} \dot{X}(t, \mathbf{x}) = b(X(t, \mathbf{x})) \\ X(0, \mathbf{x}) = \mathbf{x} \\ X\left(\frac{1}{2}, \mathbf{x}\right) = I_{-}\left(\psi(I_{+}^{-1}(\mathbf{x}))\right) \end{cases}$$
(3.2)

It is easy to see that (3.2) admits a unique solution X_{ψ} . We have to prove that X_{ψ} is measure preserving. A computation like (3.1) shows that $X_{\psi}(t, \mathbf{E}) \# \mathcal{L}^3 = \mathcal{L}^3(\mathbf{E})$ for all sets \mathbf{E} of the form

$$\mathbf{E} = \{ (r, \theta, z) : \theta \in E_1, r \in [0, \sqrt{z}], z \in E_2 \}$$
 (3.3)

where $E_1 \subset [0, 2\pi], E_2 \subset \mathbb{R}$. Sets of the form (3.3) are a basis for the Borel σ -algebra, hence X_{ψ} preserve the 3D Lebesgue measure on Borel sets. Since X_{ψ} maps null sets into null sets, it follows that it is a regular Lagrangian flow.

4. Some ideas and heuristics on possible extensions

Consider the maps

$$\psi_1(\theta) = \begin{cases} \theta & \text{if } \theta \in [0, \pi), \\ 3\pi - \theta & \text{if } \theta \in [\pi, 2\pi], \end{cases}$$

and

$$\psi_2(\theta) = \begin{cases} 2\theta & \text{if } \theta \in [0, \pi), \\ 2(\theta - \pi) & \text{if } \theta \in [\pi, 2\pi]. \end{cases}$$

The map ψ_1 leaves half a circle fixed and flips the other half, while the map ψ_2 rotates twice around \mathbb{S}^1 . Since the strategy of the proof of Theorem 2.1 produces in the limit only solutions of the form X^{Θ} , we wonder if it is possible to obtain, as limit of a suitable approximation, the flows X_{ψ_1}, X_{ψ_2} associated to ψ_1, ψ_2 as in the proof of Proposition 3.2. This is a concrete example of the following general question:

(Q2) Does the approximation procedure obtained by smoothing the vector field select a subset of the flows of b?

The strategy of [4] selects the regular Lagrangian flows corresponding to measure preserving map of the form $\psi(\theta) = \theta + \Theta \mod 2\pi$. These flows are in a sense "better" than the others for the following reasons:

- the flows X^{Θ} self intersect only in the origin, while this is not true for X_{ψ_2} , which is not even a.e. invertible;
- the Jacobian of X^{Θ} does not change sign, while this is the case for X_{ψ_1} .

Consider a general smooth approximation b_{ε} of the vector field b; the corresponding Cauchy problem admits a uniquely defined sequence of flows X^{ε} and one can ask to which X_{ψ} the sequence X^{ε} may converge. It is not clear to us if it is possible to construct an approximation of b in such a way that the approximated flow converge to X_{ψ_1} or X_{ψ_2} , especially if we want to approximate b only close to the singularity at the origin. We can however provide some heuristics motivating why it is not trivial to exclude the possibility of getting X_{ψ_1} in the limit just by arguing on the base of "topological obstructions". In fact, we can approximate the flow X_{ψ_1} with maps X^{ε} of the form:

$$X^{\varepsilon}(t,\mathbf{x}) = \begin{cases} X(t,\mathbf{x}) & \text{for } 0 \leq t \leq t_{1}^{\varepsilon} := \frac{z^{2} - \varepsilon^{2}}{4}, \\ \frac{t - t_{1}^{\varepsilon}}{t_{2}^{\varepsilon} - t_{1}^{\varepsilon}} I_{-}\left(\psi_{1}(I_{+}^{-1}(\mathbf{x}))\right) + \frac{t_{2}^{\varepsilon} - t}{t_{2}^{\varepsilon} - t_{1}^{\varepsilon}} X(t_{1}^{\varepsilon},\mathbf{x}) & \text{for } t_{1}^{\varepsilon} \leq t \leq t_{2}^{\varepsilon} := \frac{z^{2}}{4} + \frac{\varepsilon^{2}}{4}, \\ X\left(t - t_{2}^{\varepsilon}, I_{-}\left(\psi_{1}(I_{+}^{-1}(\mathbf{x}))\right)\right) & \text{for } t_{2}^{\varepsilon} \leq t < \infty, \end{cases}$$

where $\mathbf{x} \in P^+$. Each X^{ε} is a well-defined map, which is however not a flow a vector field. Therefore, this does not answer our question. However, this example tells us that an answer in the positive to our question could not just rely on topological properties of the approximating flows.

Acknowledgments. This research has been supported by the ERC Starting Grant 676675 FLIRT.

References

- [1] L. Ambrosio: Transport equation and Cauchy problem for BV vector fields. Inventiones Mathematicae, 158, (2004), 227-260.
- [2] L. Ambrosio, G. Crippa: Continuity equations and ODE flows with non-smooth velocities. Proc. Roy. Soc. Edinburgh Sect. A, 144, (2014), 1191-1244.
- [3] F. BOUCHUT, G. CRIPPA: Lagrangian flows for vector fileds with gradient given by a singular integral. J. Hyperbolic Diff. Equ., 10, (2013), 235-282.
- [4] G. CIAMPA, G. CRIPPA, S. SPIRITO: Smooth approximation is not a selection principle for the transport equation. Preprint.
- [5] G. CRIPPA, C. DE LELLIS: Estimates and regularity results for the DiPerna-Lions flow. J. Reine Angew. Math., 616, (2008), 15-46.
- [6] G. CRIPPA, C. NOBILI, C. SEIS, S. SPIRITO: Eulerian and Lagrangian solutions to the continuity and Euler equations with L¹ vorticity. SIAM J. Math. Anal. 49 (2017), no. 5, 3973–3998.
- [7] N. Depauw: Non unicit des solutions bornes pour un champ de vecteurs BV en dehors d'un hyperplan. C.R. Math. Sci. Acad. Paris, 337, (2003), 249-252.
- [8] R.J. DIPERNA, P.L. LIONS: Ordinary differential equations, transport theory and Sobolev spaces. Invententiones Mathematicae, 98, (1989), 511-517.
- [9] Q. H. NGUYEN: Quantitative estimates for regular Lagrangian flows with BV vector fields. http://cvgmt.sns.it/media/doc/paper/3848
- (G. Ciampa) GSSI GRAN SASSO SCIENCE INSTITUTE, VIALE FRANCESCO CRISPI 7, 67100 L'AQUILA, ITALY & DEPARTMENT MATHEMATIK UND INFORMATIK, UNIVERSITÄT BASEL, SPIEGELGASSE 1, CH-4051 BASEL, SWITZERLAND

E-mail address : gennaro.ciampa@gssi.it,gennaro.ciampa@unibas.ch

(G. Crippa) Department Mathematik Und Informatik, Universität Basel, Spiegelgasse 1, CH-4051 Basel, Switzerland

 $\textit{E-mail address} : {\tt gianluca.crippa@unibas.ch}$

(S. Spirito) DISIM - DIPARTIMENTO DI INGEGNERIA E SCIENZE DELL'INFORMAZIONE E MATEMATICA, UNIVERSITÁ DEGLI STUDI DELL'AQUILA, VIA VETOIO, 67100 L'AQUILA, ITALY

 $E\text{-}mail\ address: \verb|stefano.spirito@univaq.it||$

LATEST PREPRINTS

No.	Author: Title
2017-10	M. J. Grote and U. Nahum Adaptive Eigenspace Regularization For Inverse Scattering Problems
2017-11	J. Dölz, H. Harbrecht, S. Kurz, S. Schöps and F. Wolf A Fast Isogeometric BEM for the Three Dimensional Laplace- and Helmholtz Problems
2017-12	P. Zaspel Algorithmic patterns for H-matrices on many-core processors
2017-13	R. Brügger, R. Croce and H. Harbrecht Solving a free boundary problem with non-constant coefficients
2017-14	M. Dambrine, H. Harbrecht and B. Puig Incorporating knowledge on the measurement noise in electrical impedance tomography
2017-15	C. Bürli, H. Harbrecht, P. Odermatt, S. Sayasone and N. Chitnis Analysis of Interventions against the Liver Fluke, Opisthorchis viverrini
2017-16	D. W. Masser Abcological anecdotes
2017-17	P. Corvaja, D. W. Masser and U. Zannier Torsion hypersurfaces on abelian schemes and Betti coordinates
2017-18	F. Caubet, M. Dambrine and H. Harbrecht A Newton method for the data completion problem and application to obstacle detection in Electrical Impedance Tomography
2018-01	H. Harbrecht and P. Zaspel On the algebraic construction of sparse multilevel approximations of elliptic tensor product problems
2018-02	F. Ghiraldin and X. Lamy Optimal Besov differentiability for entropy solutions of the eikonal equation
2018-03	H. Harbrecht and M. Schmidlin Multilevel quadrature for elliptic problems on random domains by the coupling of FEM and BEM
2018-04	M. Bugeanu and H. Harbrecht Parametric representation of molecular surfaces
2018-05	A. Abdulle, M. J. Grote and O. Jecker Finite element heterogeneous multiscale method for Elastic Waves in Heterogeneous Media

LATEST PREPRINTS

No.	Author: Title
2018-06	M. J. Grote and J. H. Tang On controllability methods for the Helmholtz equation
2018-07	H. Harbrecht and M. Moor Wavelet boundary element methods — Adaptivity and goal-oriented error estimation
2018-08	P. Balazs and H. Harbrecht Frames for the solution of operator equations in Hilbert spaces with fixed dual pairing
2018-09	R. Brügger, R. Croce and H. Harbrecht Solving a Bernoulli type free boundary problem with random diffusion
2018-10	J. Dölz, H. Harbrecht and M. D. Multerer On the best approximation of the hierarchical matrix product
2018-11	H. Harbrecht and P. Zaspel A scalable H-matrix approach for the solution of boundary integral equations on multi-GPU clusters
2018-12	H. Harbrecht, N. Ilić and M. D. Multerer Acoustic scattering in case of random obstacles
2018-13	D. H. Baffet, M. J. Grote, S. Imperiale and M. Kachanovska Energy decay and stability of a perfectly matched layer for the wave equation
2018-14	D. Baffet and M. J. Grote On wave splitting, source separation and echo removal with absorbing boundary conditions
2019-01	M. Graff, M. J. Grote, F. Nataf and F. Assous How to solve inverse scattering problems without knowing the source term: a three-step strategy
2019-02	Z. Gao and P. Habegger Heights in families of abelian varieties and the geometric Bogomolov conjecture
2019-03	M. J. Grote, F. Nataf, J. H. Tang and PH. Tournier Parallel Controllability Methods For the Helmholtz Equation
2019-04	G. Ciampa, G. Crippa and S. Spirito On smooth approximations of rough vector fields and the selection of flows