

Supplementary Information

Induced Fit and Mobility of Cycloalkanes within Nanometer-sized Confinements at 5K

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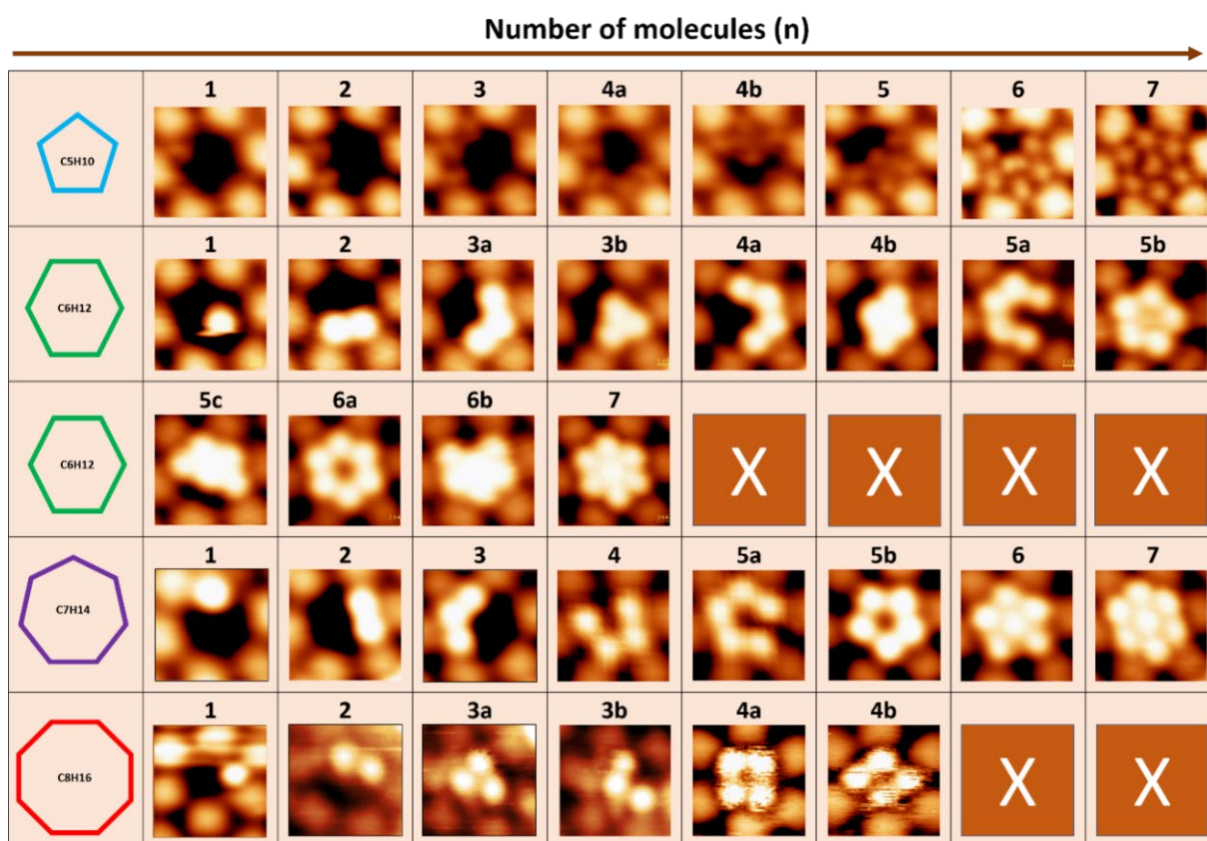
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48 **1. Molecule-by-molecule condensation of c-alkanes.**

49 Samples covered in part by DPDI networks have been exposed to different cyclo-alkanes and
50 investigated by STM. The complete data set showing all observed configurations of **c5** to **c8** in
51 the pores is shown here (Figure S1.1). The cyclo-octane molecules in the situation 4a in Figure
52 S1.1 are mobile within the pore and their average residence time has been estimated from the
53 streak dimensions in cross sections of the STM data. From the timing of the scanlines and the
54 length of the streaks, we estimate an average residence time of the order of ~60 msec. The **c8**
55 case is extraordinary in multiple aspects, (i) as the threefold symmetry of the network is broken
56 by the asymmetric triangle / rectangular arrangement formed at occ(3) and occ(4) respectively
57 and also (ii) in the considerable mobility of the molecules as indicated by the noise in the
58 scanlines. This level of mobility is unexpected at 4K and specifically occurs in the pore, not
59 for c-octanes adsorbed on the metal terrace.



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61 **Figure S1: Molecule by molecule condensation of c-alkanes inside confinements.** STM images
 62 (2.4nm x 2.4nm) of pores containing different number of c-alkane molecules. The pictogram at the left
 63 side labels the concerned c-alkane and the number of molecules / pore is indicated by the numbers (n:1-
 64 7) above each image where 'a' and 'b' denote condensates comprising the same number of molecules
 65 in different configurations as identified from time lapse images. STM parameters: 1V, 6pA, image size:
 66 2.4 nm x 2.4 nm, 256 pixels per frame, scan speed 2nm/sec.

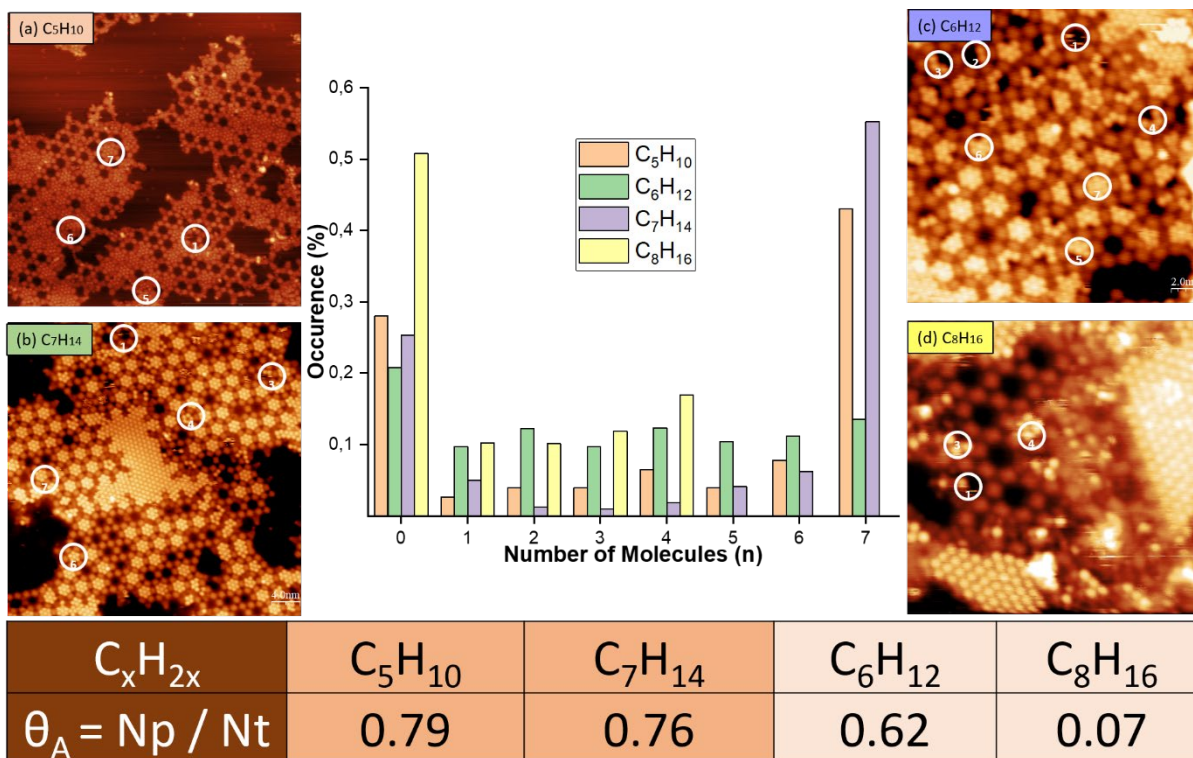
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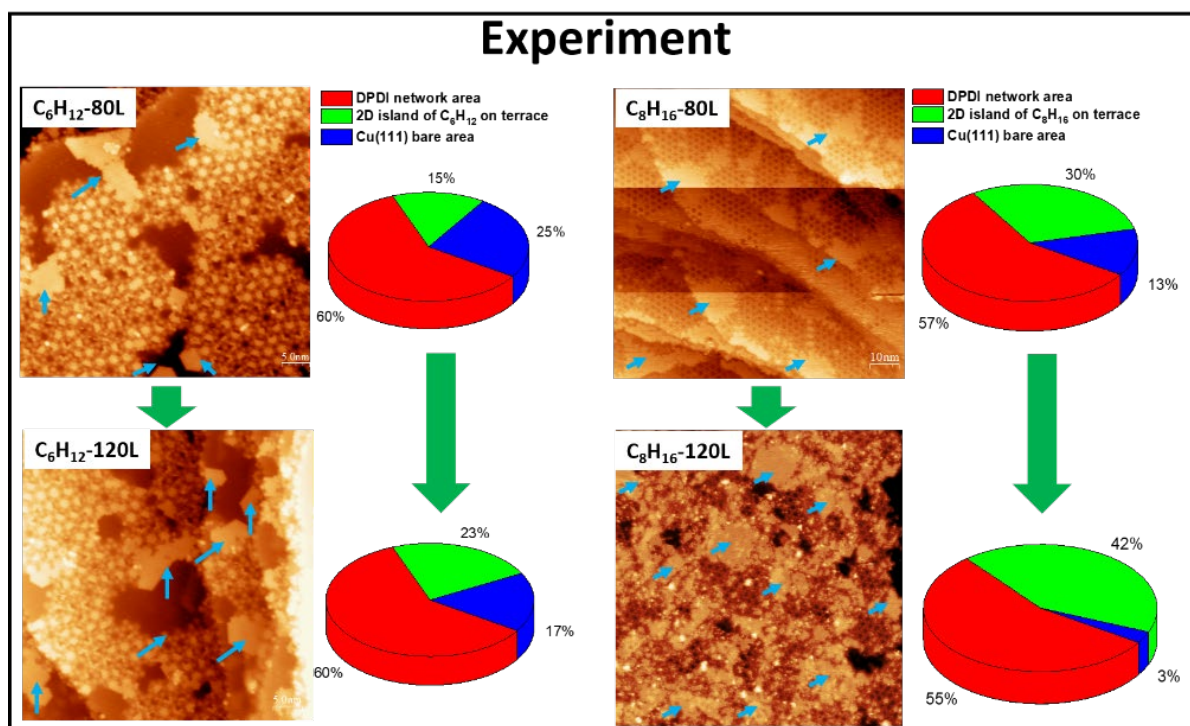
69 2. Molecular adsorption / condensation statistics

70 To assess the adsorption / condensation statistics, the relative adsorption in 'network pores' vs
 71 the total adsorption on the surface has been tabulated at the bottom of Figure S2. In analogy to
 72 the *Langmuir Adsorption Model*, we can take this as an evidence that the adsorption probability
 73 for c-alkanes to adsorb in the DPDI pores is significantly higher for **c5** & **c7** (0.79 & 0.76) in
 74 comparison to **c6** (0.62). Most remarkably, **c8** shows the least tendency of adsorbing inside the
 75 pores (0.07). This non-linear behaviour with the order number and size of the c-alkane is
 76 related to the combined effect of spatial confinement and interaction with the quantum well
 77 state as explained in the main manuscript.

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80 **Figure S2.1: Histogram analysis of c-alkanes inside DPDI confinements.** (a – d) STM micrographs
81 indicating the considerably different adsorption statistics of the investigated c-alkanes in the DPDI
82 network. All samples have been held at 9K while being exposed to 120L c-alkane gas. After exposure,
83 samples have been annealed to 20K and imaged by STM. Around 600 pores have been analysed from
84 different large scale STM micrographs for each of the c-alkanes. Before exposure, all samples have
85 been covered by DPDI network islands covering ~70 % of the surface area. Molecules inside the pores
86 and on the free metal surface can be discriminated by their position with respect to the darker network
87 backbone. The statistical distribution of filled and empty pores strongly depends on the particular
88 cycloalkane and is presented in the centre of the figure as histogram plot. The fractional occupancy (θ_A=
89 N_p / N_t) is tabulated at the bottom of the figure (N_t: total number of molecules counted; N_p: number of
90 molecules adsorbed in the pores). The fractional occupancy θ_A is considerably lower for the largest c-
91 alkane i.e. **c8** than for the other investigated c-alkanes.



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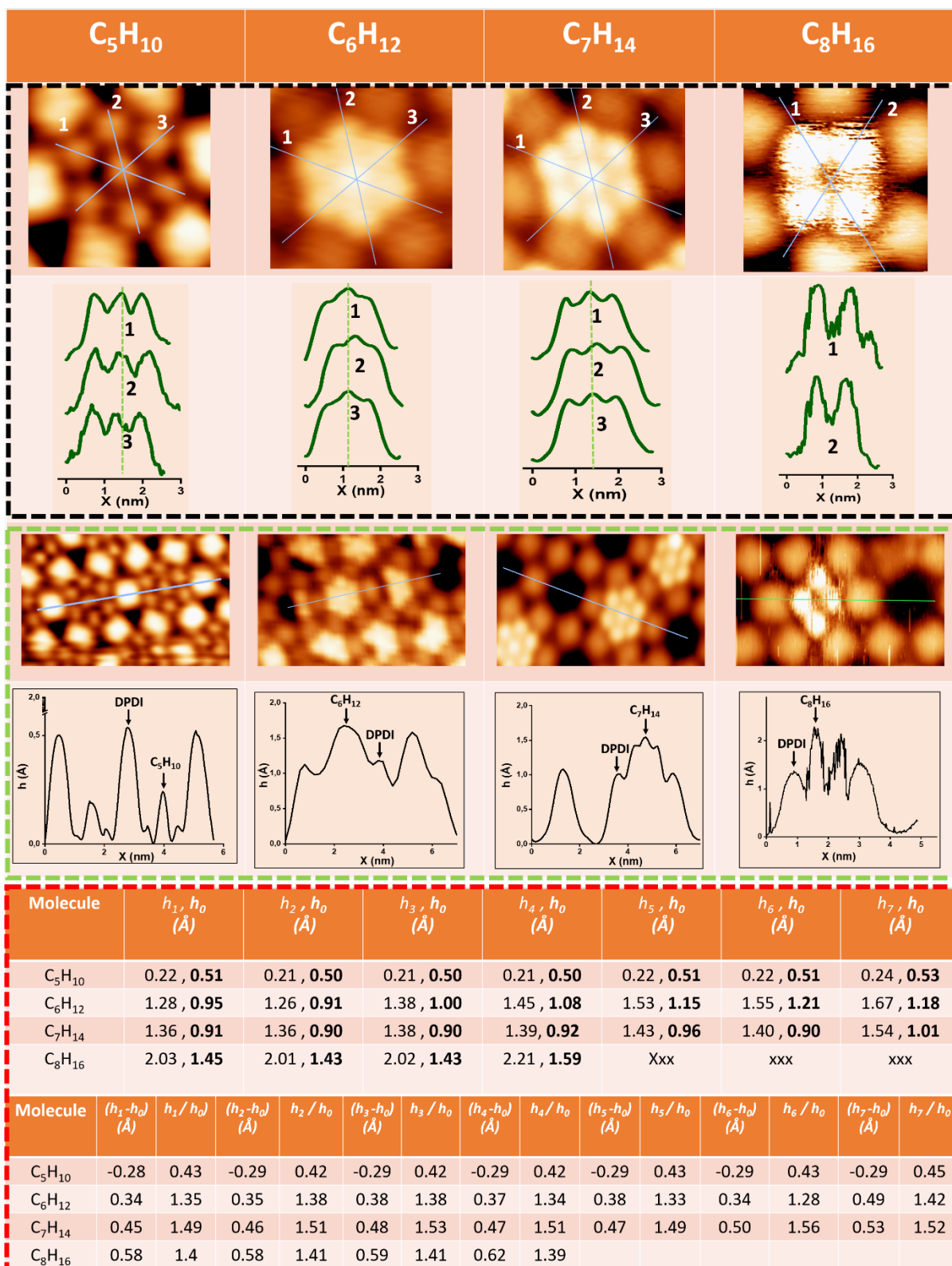
94 **Figure S2.2: Coverage dependent distribution of c-alkanes on network covered and**
 95 **network free area.** Left hand side: surface coverage after 80L and 120L exposure of **c6**. The
 96 molecular coverage on the network-free, bare Cu(111) terrace increases gradually as exposure
 97 time increases. In the STM micrographs for **c6**, a large percentage of the pores is at least
 98 partially occupied, even after the lower exposure. Right hand side: surface coverage at 80L and
 99 120L exposure of **c8**. Interestingly the surface coverage of **c8** on bare Cu(111) terraces is
 100 comparably high and accompanied by a low fraction of at least partially occupied pores of the
 101 network. This data set and histogram analysis indicate that the sticking / adsorption events in
 102 the DPDI pores are reduced with increased molecular dimensions. During the deposition
 103 process, **c8** relocates from network covered areas to the free Cu surface. Note that the fraction
 104 of DPDI network on the samples has been held approximately constant to allow for a valid
 105 comparison. Blue arrows in the micrographs indicate the bare Cu(111) surface fraction.

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3. Height analysis of c-alkanes inside pores and on bare Cu

109 Different apparent height analyses of c-alkanes inside confinements are provided in
 110 comparison to their apparent height upon adsorption on the network free area are presented in
 111 Figure S3.1. The apparent height is shown to depend on both, the position and arrangement of
 112 c-alkanes in individual pores (Figure S3.1, top panel outlined in black) and can be referenced
 113 to the apparent height of the DPDI network backbone (Figure S3.1, center panel outlined in
 114 green). Also, there is a change in the apparent height depending on whether the molecules are
 115 adsorbed inside the pores or on the network free terraces (Figure S3.2).

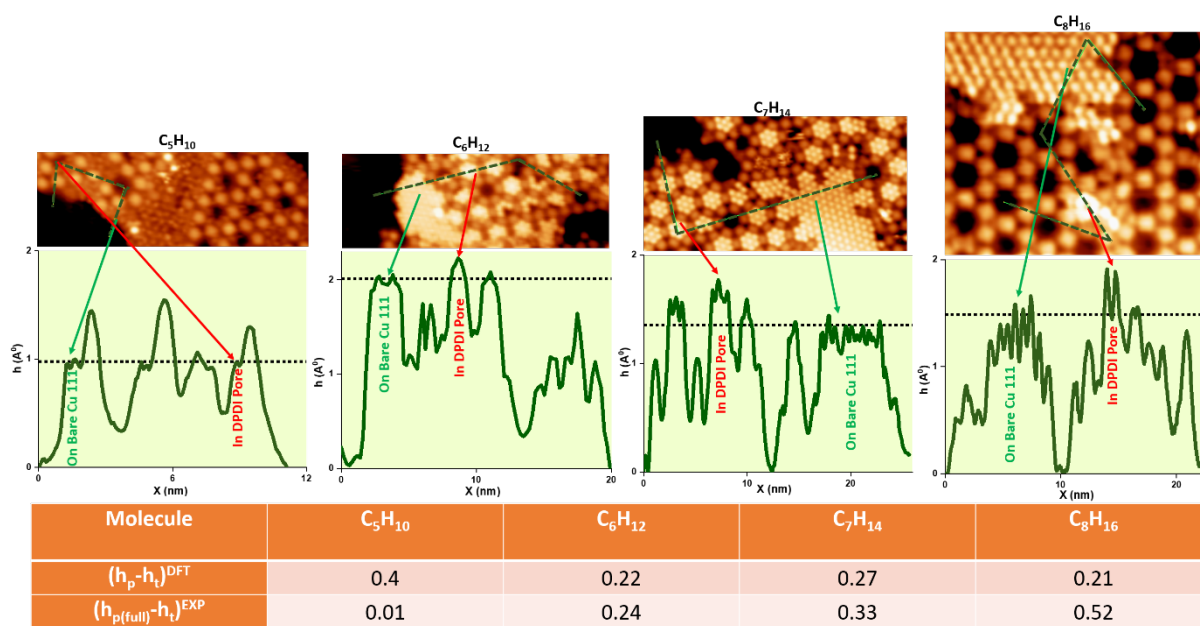
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119 **Figure S3.1: Apparent STM height: profile analysis of c-alkanes in the DPDI confinements.**
 120 Three panels have been compiled with STM data and simulations which are: i) position dependent
 121 apparent STM height of c-alkanes inside pores (*black outlined table*), ii) the relative height of c-alkanes
 122 with respect to the DPDI network backbone (*green outlined table*) and iii) the occupancy dependent
 123 evolution of the apparent height (*red outlined table*). *Black table*: STM micrographs of the completely
 124 c-alkane-filled pores are shown at the top. The in-pore intermolecular distances between the individual

125 molecules and the apparent STM height can be analysed from the cross sections shown below the STM
 126 micrographs. Note that the center molecule for **c6** and **c7** in the fully occupied pores is protruding
 127 slightly higher than its six neighbours. This is an effect which we tentatively associate with steric
 128 constraints upon full packing of the pore. *Green table:* The apparent height of c-alkanes within the pores
 129 in comparison to the DPDI network backbone is characteristically different for **c5** and **c6-c8** as it can
 130 be recognized in the STM micrographs and cross sections. *Red table:* All determined height parameters
 131 have been provided and compared as they are discussed in the main manuscript; absolute and relative
 132 comparisons are listed. The difference between the height of the molecules at occupancy $occ(i)$ is
 133 tabulated as h_i and the height of the DPDI is tabulate as h_0 . All values are given in Å. STM parameters:
 134 1V, 6pA, pixels per frame: 256 px, scan speed 2nm/sec.
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 138 **Figure S3.2: Apparent STM height profile analysis of c-alkanes in fully filled DPDI pores in**
 139 **comparison with supramolecular islands on bare Cu(111).** The STM apparent height of the fully
 140 occupied pore, $h_{p(full)}$ has been compared to the corresponding height of extended 2D islands on
 141 neighbouring terraces h_t . The apparent height differences increase monotonously from **c5** to **c8**
 142 indicating an increasingly strong site-specific interaction inside the pore. This trend is well captured in
 143 the DFT calculations (row labelled with $(h_p - h_t)^{DFT}$ in the table at the bottom of the figure) with the
 144 increase towards **c8** being lower than experimentally observed. We tentatively assign this exception to
 145 the experimentally observed dynamicity of **c8** on the time scale of ~60 msec which may not be captured
 146 by DFT in a straightforward way. Note that the data in the table has been obtained from a small number
 147 (~3-5) of representative cross sections taken at different places across the sample. STM parameters: 1V,
 148 6pA.
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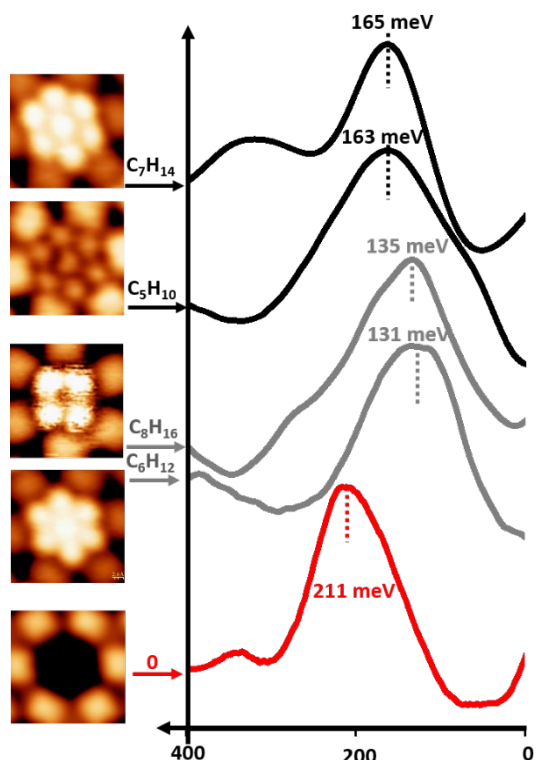
150 **c5** exhibits an exceptionally low apparent height of ~ 43% of the apparent height of the DPDI
 151 backbone (Figure S3.1). The close contact between **c5** at full ($occ=7$) occupancy of the pore
 152 may also cause the interesting trapezoidal shape of the center molecule in the STM micrographs
 153 (Figure S3.1 top right). In the experiment, the higher order number c-alkanes **c6 – c8** exhibit a
 154 considerably larger apparent height, higher than the DPDI backbone (Further details of the
 155 apparent height analysis are tabulated in Figure S3.1 & Figure S3.2).
 156

157 **4. DFT Estimate of binding energy of c-alkanes in pores**

Conformer	ΔE_0	ΔE_p^1	$\Delta E_p^2/2$
C_5H_{10}	-0.47	-0.67	-0.57
C_6H_{12}	-0.49	-0.67	-0.59
C_7H_{14}	-0.61	-0.74	-0.70
C_8H_{16}	-0.66	-0.79	-0.74

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 160 **Figure S4: Binding energy of cyclo-alkanes in pores: decreasing energy per molecule with**
 161 **occupancy increasing from 1 to 2.** The binding energy is given in eV; ΔE_0 is binding energy on bare
 162 surface while ΔE_p^n ($n = 1, 2$) is the binding energy inside pore for single ($n=1$) or two ($n=2$) molecules.
 163 Note that DFT only captures the Energy ground state which in conjunction with the observed
 164 phenomena may not be the relevant state, in particular at higher packing density.

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 167 **5. Pauli repulsion between confined surface state and the c-alkanes**



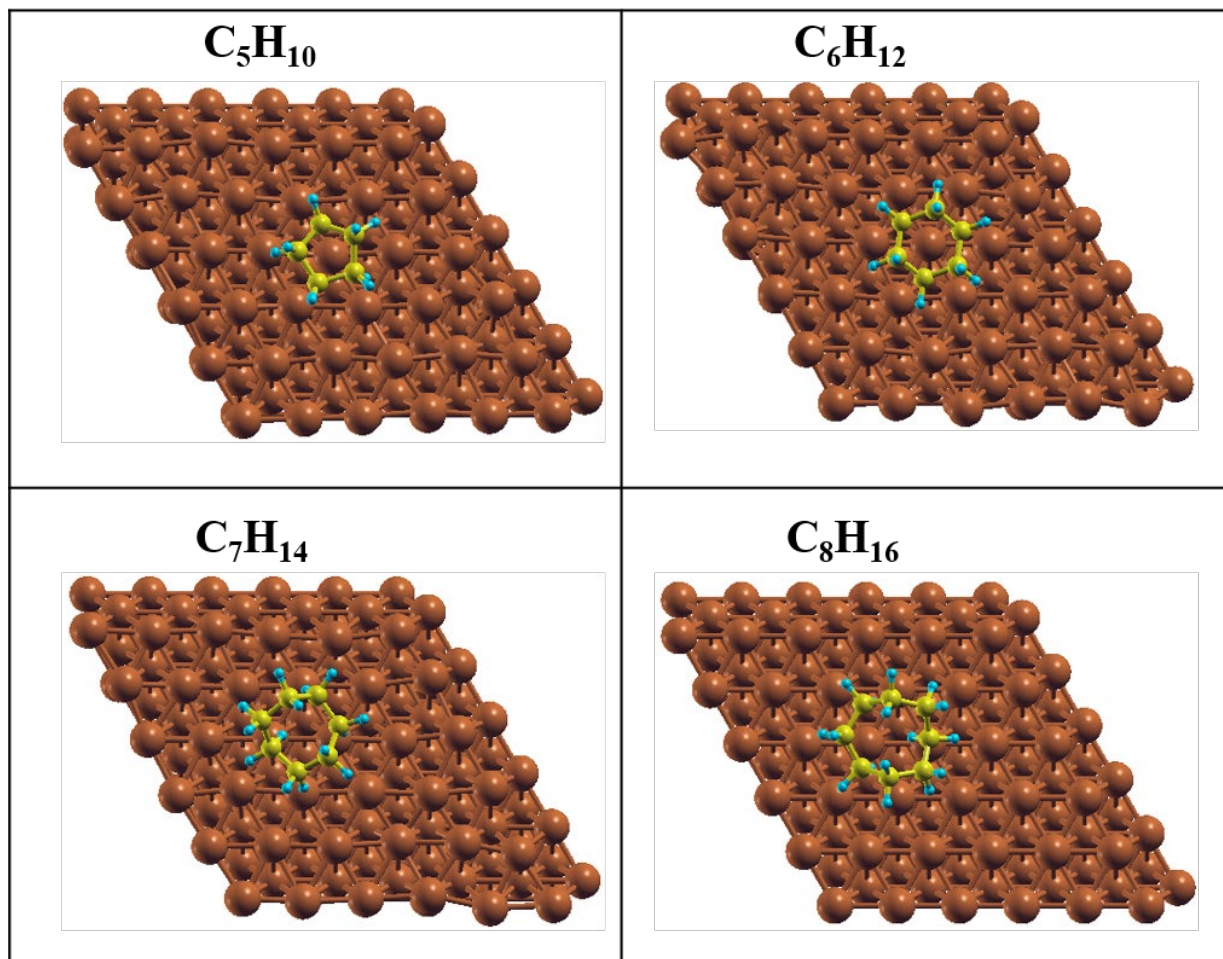
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 169 **Figure S5: Pauli Repulsion: Fully filled c-alkanes interacting with the confined surface state.**
 170 Comparing the Pauli repulsion, the c-alkanes containing an uneven number of C-atoms (**c5** and **c7**)
 171 exhibit a binding energy (BE) shift which is ~ 30 meV closer to the BE of the native confined surface

172 state (CSS) (lower Pauli repulsion, more relaxed) than their even order number counterparts **c6** and **c8**.
173 This is taken as a sign for the changing adsorption mode as discussed in the main paper.

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175 6. C-alkane conformation on Cu(111)

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179 **Figure S6: c-alkanes after relaxation on Cu(111).** 3D plots of the atomic structures obtained with the
180 DFT calculations presented in Table 1 of the main paper. *Top-left:* the lowest energy configuration of
181 C_5H_{10} on Cu(111) is close to the envelope configuration which is degenerate in that each C-atom can
182 take the role of the ‘flap’. *Bottom-left:* C_7H_{14} assumes a less close to planar configuration suggesting
183 that the stabilization via the surface is less effective. *Second column:* C_6H_{12} adsorbs in an only slightly
184 closer to planar chair configuration and C_8H_{16} assumes a complex shape which is again less coplanar
185 with the surface. The decreased planarity of molecular module facilitates a lower adsorption energy and
186 residual dynamicity compared to planar or strongly interacting molecules. (see also Table 1 in main
187 paper)

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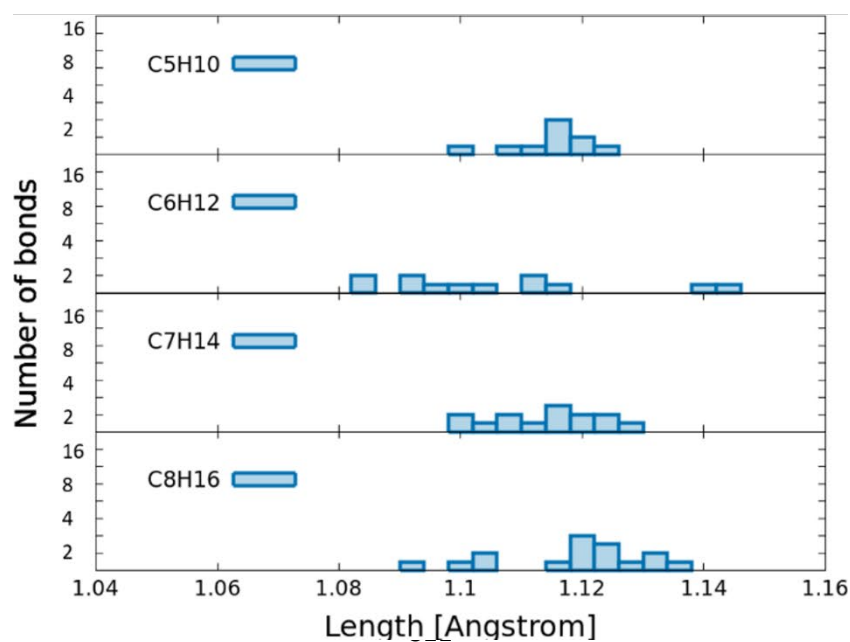
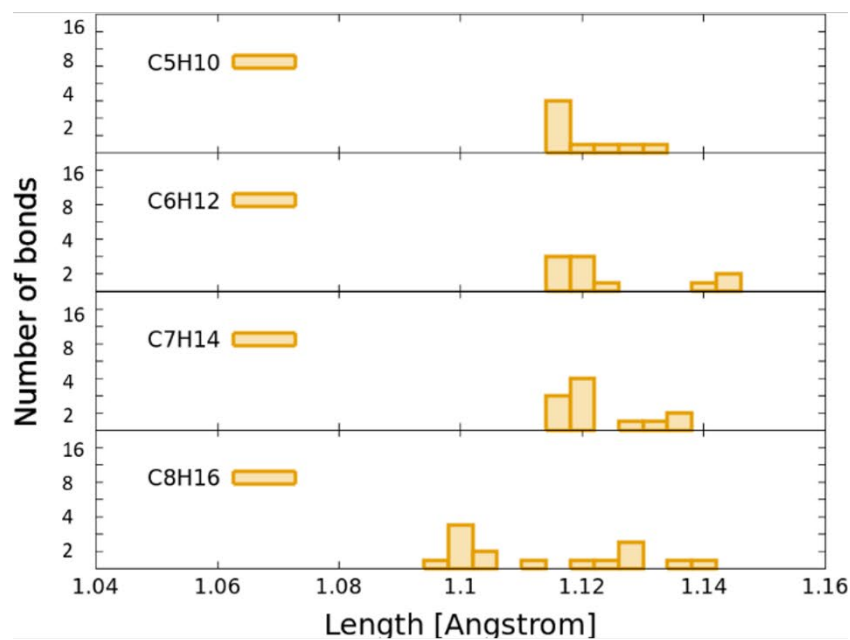
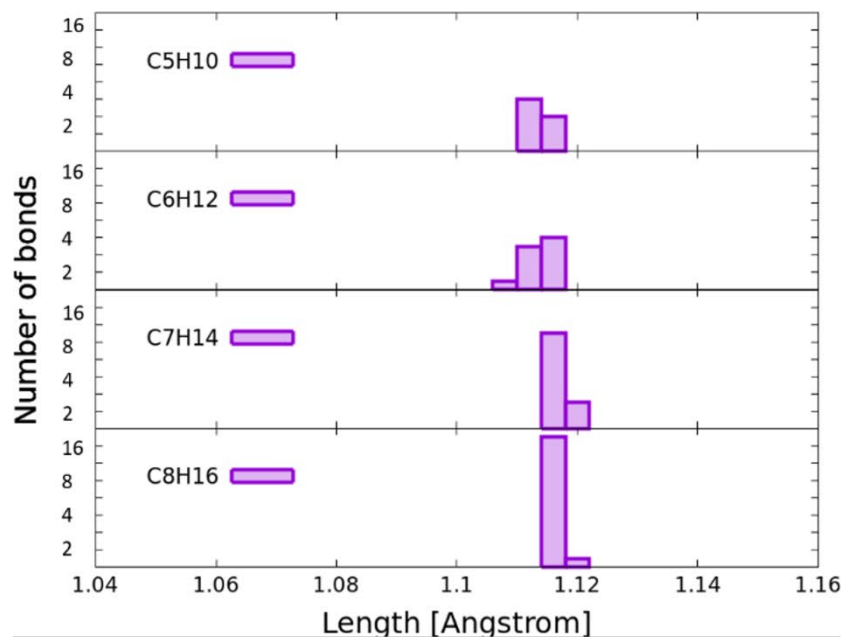
191 7. Bond length comparison of c-alkanes.

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193 For free molecules (i.e. no influence of the surface) we see that almost all bonds share the same
194 length, close to 1.12 Angstroms as shown in (Figure S-7). For the molecules adsorbed on clean
195 Cu(111) surface (yellow plots) **c5**, **c6** and **c7** have a clear tendency to increase their C-H bond
196 lengths (i.e. the C-H bond is stretched by the forces acting in the surface potential). For **c8**, we
197 note both shorter and longer C-H bonds. This indicates, that the interaction of **c8** with the
198 surface is different from the other three cases. Finally, upon adsorption of the c-alkanes in the
199 pore, the situation becomes much more complex, with shorter C-H bonds for **c5** and mixtures
200 of shorter/longer bonds for the other molecules. Nevertheless, the plot suggests a trend to form
201 shorter C-H, indicating that its interaction with the surface is different from the other three.
202 While the changes in the C-H lengths are relatively small, the trends are sufficiently clear to
203 conclude that the interaction between molecules and clean surface is different from that
204 between molecule and the surface state inside the pore. Indeed, the surface state leads to a shift
205 of the Fermi level which in turn influences the molecule-surface charge transfer / interaction
206 mechanisms. This has two reasons; (1) changes in the line-up of the Fermi level, and (2) the
207 position of molecular HOMO/LUMO. An inspection of the aforementioned data shows that
208 this effect is very weak for **c8** while it is maximum for the **c6**. Note that this simulation has
209 only been performed for a few molecules inside the pores. Therefore, it is indicative of the
210 complexity of the mechanisms involved in the process, but not a realistic representation of the
211 experimental situation. The observed changes in the simulation are strongly suggesting that
212 upon pore filling similar bond-length and bond angle, modifications occur as they are expected
213 to cause the phenomena as they are explained by the ‘induced fit’ concept in the main paper.

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215



217 **Figure S7. Graphical representation of the C-H bond-lengths** for the free molecules (purple),
218 modified by their adsorption in supramolecular islands on Cu(111) (yellow) and by their adsorption
219 inside the pore (blue). The histogram nicely reflects the uniform bond length for the minimal energy
220 conformation in vacuum, the bond length spreading by the molecule-surface interaction in the contact
221 zone leading to unequal bond extension. In the third case the Pauli interaction with the confined surface
222 state ‘softens’ the hard sphere interaction of the contact atoms with the adsorbate and therefore an in-
223 between position is assumed between the in-vacuum (relaxed) and the on-surface (bottom side
224 supported) conformations and bond length spectra. Each bar shows the number of C-H bonds with a
225 certain length, starting from the value of 1.04 Angstroms, with a step size of 0.03 Angstroms.

226

227 **8. Methods**

228

229 **Sample preparation and STM measurement:**

230 The samples have been prepared and examined in an ultrahigh vacuum (UHV) system with a
231 base pressure of 6×10^{-11} mbar. The Cu(111) crystal (MaTeck GmbH) has been prepared by
232 rounds of Ar^+ sputtering at $E = 1$ keV performed at room temperature followed by annealing at
233 480°C . The DPDI molecules have been deposited from a nine-cell commercial evaporator
234 (Kentax, GmbH, Germany) on the Cu(111) by sublimation at $\sim 240^\circ\text{C}$ and the rate has been
235 controlled before deposition by a quartz crystal microbalance. After deposition, the sample has
236 been annealed to $290\text{--}300^\circ\text{C}$ in order to convert DPDI into 3deh-DPDI, which crafts the Cu-
237 coordinated network. Cycloalkanes of purity 99% has been cleaned from air contamination by
238 freeze – pump – thaw – cycles. A leak valve equipped with a capillary has been used to direct
239 the molecular vapour closer (~ 20 cm) to the sample in the STM setup. Thus, contamination of
240 the cryogenic system shields by the dosing experiments has been reduced. During the dosing
241 procedure, the STM tip had been visibly retracted (~ 100 micrometer) from the surface. The
242 STM has been operated at the equilibration temperature of the He bath cryostat (4.2K). By
243 opening of the cryo-shields during the molecular deposition process, the T reference
244 measurement increased up to ~ 5.5 K. All STM data presented in Figure1-Figure 4 (main paper)
245 and in Figure S1-S5 have been acquired after exposure of the Cu-coordinated 3deh-DPDI
246 network to 120 L to the corresponding cycloalkanes. During the exposure an overall chamber
247 pressure equal to 1.3×10^{-7} mbar has been measured. 1,200 s was the typical during exposure
248 time for the samples held at ~ 9 K. After exposure with corresponding c-alkanes, samples were
249 annealed upto 20K. The STM was operated after pumping down to the base pressure and with
250 closed cryoshields in the He cryostat ($\sim 4.2\text{K}$).

251 STM measurements have been performed in the constant current mode with Pt–Ir tips (90% Pt,
252 10% Ir). The tip has been prepared by mechanical cutting followed by sputtering and controlled
253 indentation into bare Cu(111) surfaces. If not specified differently, 1V sample bias has been
254 applied and the feedback circuit has been set to a tunnelling current of 6pA.

255

256 **DFT computational setup**

257 Simulations were performed using the SIESTA code^{1,2}. As exchange-correlation functional,
258 we have used the vdW-DF-cx functional of Berland and Hyldgaard (BH)³. The advantage of

259 this functional over other van der Waals functionals is that it accurately reproduces the structure
260 of the metallic bulk. By calculating crystalline Cu we obtain a lattice constant of 3.64 Å; this
261 comprises a value 0.8% difference than experimental value (3.61 Å). As basis sets, we used
262 double-zeta polarized (DZP) with an energy shift of 50 meV for Cu slab, while for the
263 cycloalkanes we have used triple-zeta polarized basis sets (TZP) with an energy shift of 25
264 meV.

265 The geometry used for the surface decorated with DPDI pores is that reported by Matena et al;
266 ⁴ to simulate the copper surface we used 4 atomic layers, passivated by a layer of H and
267 accommodated into a 10x10 supercell on the Cu(111) surface. 6x6 super-cell dimensions have
268 been chosen for calculations on the bare surface in absence of the network. The length of the
269 cells along the OZ axis was 40 Å for all systems, thus permitting a vacuum level of ≈ 30 Å and
270 allowing us to avoid the artificial influence of the electric charge from one cell to another.

271 For all the systems, we calculated Gamma point integrals in the Brillouin zone. Structural
272 relaxation was undertaken in two steps; in order to assess the molecule's intrinsic flexibility,
273 we have used the 'simulated annealing' method with an initial temperature T=500K and a final
274 annealing temperature of 2K. The time step of the dynamics in the sequences was 1 fs. Next,
275 the systems were minimized until the maximum gradient in the relaxed structure was below
276 0.01 eV/Å by using the Conjugated Gradient algorithm.

277 Simulated STM images have been obtained by using the Tersoff-Hamman approximation⁵ as
278 implemented by Lorente and Persson⁶. We computed the STM images for a bias of 1 V in a
279 simulation of a 'constant current' STM experiment (i.e. following a constant value of local
280 density of states in Tersoff Hamman approximation).

281 We have performed structural relaxation for all relevant conformers reported in the literature
282 in order to assess the validity of our methods. Our results are in accord to those already reported
283 [c5-c8]⁷⁻¹⁰.

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