

Optimal Besov differentiability for entropy solutions of the eikonal equation

F. Ghiraldin and X. Lamy

Departement Mathematik und Informatik
Fachbereich Mathematik
Universität Basel
CH-4051 Basel

Preprint No. 2018-02
February 2018

www.math.unibas.ch

OPTIMAL BESOV DIFFERENTIABILITY FOR ENTROPY SOLUTIONS OF THE EIKONAL EQUATION

FRANCESCO GHIRALDIN AND XAVIER LAMY

ABSTRACT. In this paper we study the Eikonal equation in a bounded planar domain. We prove the equivalence among optimal Besov regularity, the finiteness of every entropy production and the validity of a kinetic formulation.

1. INTRODUCTION

In this paper we study the Eikonal equation in a planar domain: given a bounded Lipschitz domain $\Omega \subset \mathbb{R}^2$ we consider solutions $m : \Omega \rightarrow \mathbb{R}^2$ to the following constrained equation

$$|m| = 1 \quad \text{a.e. in } \Omega, \quad \nabla \cdot m = 0 \quad \text{in } \mathcal{D}'(\Omega). \quad (\text{M})$$

The eikonal equation (M) is very flexible, and uniqueness or regularity cannot be expected for such weak solutions, even imposing a boundary datum (the equation, on simply connected domains, is equivalent to solving $|\nabla u| = 1$ for a scalar function u , for which for instance the theory of viscosity solutions singles out a distinguished subclass of solutions).

On the other hand, solutions to (M) coming from physical models usually possess extra information that limit this flexibility. This equation emerges in the description of several physical phenomena, collectively called line-energy Ginzburg-Landau models, that describe for instance smectic liquid crystals, soft ferromagnetic films, blister formations, and broadly speaking phase transition phenomena where the order parameter is a gradient [20].

For example one can consider the Aviles-Giga energy

$$AG_\varepsilon(u_\varepsilon, \Omega) := \int_\Omega \varepsilon |\nabla^2 u_\varepsilon|^2 + \frac{1}{\varepsilon} (1 - |\nabla u_\varepsilon|^2)^2 dx, \quad \Omega \subset \mathbb{R}^d \quad (\text{AG})$$

(with appropriate boundary conditions): this energy has been introduced in [5] to study liquid crystal configurations, and in the two dimensional case was considered by Gioia and Ortiz as a model energy for the deformation of thin film blisters undergoing biaxial compression [31]. The functional (AG) can be thought of as a vectorial Modica Mortola energy, where the fields are forced to be gradients; equivalently, $\varepsilon AG_\varepsilon$ can be seen as a singular perturbation (accounting for the bending energy of the film) of the elastic energy. Competition between these two terms favors concentration along the jump discontinuities of the limit gradient ∇u , with a limit energy believed to be asymptotically

$$\frac{1}{3} \int_{\text{Jump}(\nabla u)} |\nabla^+ u - \nabla^- u|^3 d\mathcal{H}^1.$$

Solutions to (M) can be obtained from sequences (u_ε) with equibounded energy $AG_\varepsilon(u_\varepsilon) \leq E$, by setting

$$m_\varepsilon := \nabla^\perp u_\varepsilon, \quad m_\varepsilon \rightarrow m, \quad (1.1)$$

and observing that any pointwise limit m satisfies the unitary constraint $|m| = 1$ as well as the linear constraint $0 = \nabla \cdot \nabla^\perp u_\varepsilon \rightarrow \nabla \cdot m$ (see [2, 20] for the precise compactness statements).

In [29], Jin and Kohn studied the energy AG_ε and its variations (under suitable boundary conditions), and discovered that the divergences $\nabla \cdot \Sigma(\nabla u_\varepsilon)$ of suitable vectorial renormalizations of the gradient fields ∇u_ε are measures providing nontrivial asymptotic lower bounds for $AG_\varepsilon(u_\varepsilon)$. The explicit form of $\Gamma - \lim_\varepsilon AG_\varepsilon$ and its domain have been subject to intensive study, see [2, 6, 20, 12], where partial results on the $\Gamma - \lim \inf$ already conjectured in [5, 29] have been obtained.

Unit vector fields m obtained through the limit procedure (1.1) enjoy further regularity properties: they are entropy solutions. After recognizing that (M) can be interpreted as a perturbation of Burgers' equation, in [20] the parallel between these vectorial renormalizations of the eikonal equation and entropy solutions of scalar conservation laws was pushed forward, and a family ENT of entropies $\Phi: S^1 \rightarrow \mathbb{R}^2$ such that $\nabla \cdot \Phi(m)$ detect the singularities of m , has been singled out. It became therefore natural to study *entropy solutions* to (M), namely vector fields satisfying the further property that

$$\nabla \cdot \Phi(m) \in \mathcal{M}(\Omega) \quad \forall \Phi \in ENT,$$

where $\mathcal{M}(\Omega)$ is the set of finite Radon measures on Ω . As in the case of hyperbolic conservation laws, such additional informations imply further regularity and compactness of the set of solutions, in the spirit of Tartar's compensated compactness [38].

A quantitative statement of compactness, in the form of fractional differentiability was afterwards obtained, only for solutions of (M) specifically arising as limits of $\nabla^\perp u_\varepsilon$ (1.1), by Jabin and Perthame in [27]. They prove that such solutions satisfy a kinetic formulation: the equilibrium function (Maxwellian)

$$\chi(x, \xi) = \mathbb{1}_{m(x), \xi > 0},$$

defined for $\xi \in \mathbb{R}^2$, solves a transport equation of the form

$$\xi \cdot \nabla_x \chi(x, \xi) = \partial_{\xi_1} \sigma_1 + \partial_{\xi_2} \sigma_2 + \sigma_3, \quad (1.2)$$

for some locally finite measures σ_ℓ , see [27, Theorem 1.1]. Recall that in the realm of scalar conservation law, the validity of a kinetic formulation is equivalent to the finiteness of all entropy productions [33]. With the help of methods coming from velocity averaging, the authors of [27] are able to prove that such solutions possess some fractional differentiability: $m \in W_{\text{loc}}^{\frac{1}{5}-, \frac{2}{3}-}$; better Sobolev regularity $W_{\text{loc}}^{\frac{1}{3}-, \frac{2}{3}-}$ was established by the same authors in a subsequent work [28]. Examples by De Lellis and Westdickenberg [19] show that this regularity is optimal in the number of derivatives (1/3) but leave room for improving the integrability.

Similar results hold for weak solutions of Burgers' equation $\partial_t u + \partial_x \frac{1}{2} u^2 = 0$ whose entropy productions are finite measures (but may change sign). This should come as no surprise since Burgers' equation formally arises when considering solutions of (M) which are small perturbations of the constant solution $m_0 = (1, 0)$ (see e.g. the discussion in [32, p.143]). In the case of Burgers' equation, solutions with finite entropy production are shown by Golse and Perthame [23] to lie in $B_{3, \infty}^{1/3}$, which is the optimal regularity according to [19].

We wish to mention also a similar model arising in the theory of micromagnetism and studied by Rivière and Serfaty in [36, 37]. There, solutions of (M) also appear as limits of sequences with bounded energy depending on a parameter ε , and they enjoy a kinetic formulation. In that model the unit constraint is imposed already at the ε level, thus enforcing a topological restriction, while the divergence free condition is only reached in the limit, via the penalization of a nonlocal term. This feature makes the limit problem quite

different from ours (motivated by the Aviles-Giga functional): there, the field m_ε possesses an H^1 lifting $e^{i\varphi}$ (excluding vortices at the ε level - Bloch lines), that enables the use of a convenient family of entropies which control jumps of the angle φ . For this model, a quite thorough study of the rectifiability properties of entropy solutions has been carried out in [4]. Similar rectifiability properties were then obtained for the present model [16] and for higher dimensional scalar conservation laws [17]. An interesting and more sophisticated model describing (almost horizontal) micromagnetism in three dimensions, exhibiting different types of transition layers (one and two dimensional Néel walls and Bloch lines) has been considered in [1].

Another distinguished subset of solutions to (M) are the so-called zero energy states, for which the field m is again as in (1.1) with the additional property that $\lim_\varepsilon AG_\varepsilon(u_\varepsilon) = 0$. Such solutions have no entropy production: $\nabla \cdot \Phi(m) = 0$ for all $\Phi \in ENT$. This yields stronger regularity and rigidity properties, as shown by Jabin, Otto and Perthame [26]: m is locally Lipschitz outside a locally finite set of points (the *vortices*, that asymptotically carry no energy), and in any convex neighborhood of one of them (say p), it holds $m(x) = \pm \frac{(x-p)^\perp}{|x-p|}$ (see also [9] for similar results in higher dimensions). Recently Lorent and Peng [30] showed that the vanishing of only two particular entropy productions (instead of all $\Phi \in ENT$) is needed to obtain this conclusion. An indication on the minimal regularity of m needed to trigger such an improvement was further studied by De Lellis and Ignat in [15], where it is proved that if $m \in W_{loc}^{\frac{1}{3},3}$ then there is no entropy production. The $\frac{1}{3}$ differentiability exponent is somehow critical in several problems, notably the problem of energy conservation for the Euler equations (Onsager conjecture) [11, 10, 18].

In this article we prove the following (see Theorem 2.6 and Section 3 for the precise definitions):

Theorem. *Let m satisfy (M). The following three conditions are equivalent:*

- (i) m has locally finite entropy production;
- (ii) m satisfies a kinetic formulation;
- (iii) $m \in B_{3,\infty,loc}^{1/3}(\Omega)$.

This Theorem improves the previous literature in several aspects: the kinetic formulation is deduced from the mere knowledge that all entropy productions are finite, instead of the stronger requirement that m be the limit of an Aviles-Giga sequence (1.1). Whether or not the latter is *strictly* stronger is a nontrivial and, to the authors' knowledge, open question (related to the upper Γ -limit of AG_ε - what can be checked by estimating the energy of a convolution is that maps $m \in B_{2,\infty}^{1/2} \cap B_{4,\infty}^{1/4} \subsetneq B_{3,\infty}^{1/3}$ are limits of Aviles-Giga sequences, see [35]). Moreover our kinetic formulation (see (KIN) below) takes a simpler form than (1.2).

The fractional differentiability $B_{3,\infty}^{1/3}$ that we deduce from the kinetic formulation entails improved integrability compared to the previous known one [26]. As already mentioned, the corresponding result for Burgers' equation is due to Golse and Perthame [23]. Their proof relies on a kinetic formulation in which the equilibrium function χ satisfies some monotonicity assumption. This monotonicity is not present in our case, which requires substantial modification of their method.

This fractional differentiability is necessary *and sufficient* (hence optimal). Moreover our calculations also show that slightly better summability (e.g. $m \in B_{3,q}^{1/3}$ for some $q < \infty$, see § 4) already triggers the aforementioned enhanced regularity (m locally Lipschitz outside a discrete set). This criticality of $B_{3,\infty}^{1/3}$ is due to the commutator estimates employed in the

argument: similarly to the case of Euler equations, “energy conservation” for functions with slightly better differentiability properties can be proved [15, 11].

The proof of the Besov regularity from the kinetic formulation (implication (ii) \Rightarrow (iii) in the Theorem) employs an *interaction estimate* due to Varadhan [39], that was used in [23] and in [21]: as in those works we build a quantity $\Delta(x, z)$ which depends on the equilibrium function $\chi(x, \xi)$ and which controls the cubic increment $|m(x+z) - m(x)|^3$. The above-mentioned interaction estimate, together with the kinetic formulation, provides an upper bound on $\int_{\Omega} \Delta(x, z) dx$ in terms of $|z|$, hence the Besov regularity. To prove (i) \Rightarrow (ii), i.e. the validity of a kinetic formulation from the knowledge of having finite entropy production, we employ a Banach-Steinhaus argument as in [17, 8]. The other implication (iii) \Rightarrow (i) follows from a careful integration by compensation inspired by [15].

After proving the above Theorem, we explore several questions that come up naturally. As already mentioned, in the model studied by Rivière and Serfaty in [36, 37], the solutions of (M) that arise can be written as $m = e^{i\varphi}$ with some control on the lifting φ . Analogues of our entropy productions and kinetic formulation play a crucial role, and the kinetic defect measure (which is linked to the kinetic formulation) provides a *sharp* lower bound for the energy [37]. We show that the corresponding property is *not* present in our case.

Another natural question regards the set of entropies necessary to obtain the Besov regularity. Lorent and Peng prove in [30] that the vanishing of only two particular entropy productions is enough to force all the entropy productions to vanish: is the mere finiteness of these two particular entropy productions also enough to ensure the optimal regularity? We are unable to fully answer this question, but adapting some arguments in [22] we do obtain some (lower) regularity.

The article is structured as follows: after some preliminary notations in the next Section 2, in Section 3 we state and prove the main Theorem, and in Section 4 we gather some results related to the zero energy states and to the above mentioned further natural questions.

Acknowledgements. The authors thank Pierre Bochart for many enlightening discussions and his participation to the proof of Proposition 3.10. F.G. is supported by the ERC Starting Grant 676675 *FLIRT - Fluid Flows and Irregular Transport*. F.G. acknowledges the kind hospitality of Institut de Mathématique de Toulouse, where part of this paper was written.

2. NOTATIONS AND STATEMENT OF THE PROBLEM

Given $\xi, \eta \in \mathbb{R}^2$ we identify them with complex number in order to define their scalar and vector products:

$$\bar{\xi}\eta = \xi \cdot \eta + i\xi \wedge \eta.$$

Equivalently: $\xi \cdot \eta = \xi_1\eta_1 + \xi_2\eta_2$, $\xi \wedge \eta = \xi_1\eta_2 - \xi_2\eta_1$. For the sake of clarity we will not always identify unit complex numbers with rotations of the plane: in such occasions, rotations by an angle θ will be denoted by $R_{\theta} \in SO(2)$.

We will measure the smoothness of our unit fields m in the scale of Besov spaces on a domain: in order to keep the notation light, we give the definition only for the exponents we need; we refer the reader to [40] for an overview of the definitions; see also [7]. If $f : \Omega \rightarrow \mathbb{R}$ and $z \in \mathbb{R}^n$, we define $D^z f(x)$ to be the increment $f(x+z) - f(x)$ if both $x, x+z \in \Omega$, and zero otherwise.

Definition 2.1 (Besov spaces on domains, [40, Theorem 1.118]). Let Ω be a bounded Lipschitz domain and $1 \leq q \leq \infty$. A function $f : \Omega \rightarrow \mathbb{R}$ belongs to $B_{3,q}^{\frac{1}{3}}(\Omega)$ if

$$\|f\|_{L^3(\Omega)} + \left(\int_0^1 \left(t^{-\frac{1}{3}} \sup_{|z| \leq t} \|D^z f\|_{L^3(\Omega)} \right)^q \frac{dt}{t} \right)^{\frac{1}{q}} < \infty.$$

The expression on the left hand side provides a quasi-norm. Moreover we denote with $B_{3,q,\text{loc}}^{\frac{1}{3}}(\Omega) = \bigcap_{U \subset \subset \Omega} B_{3,q}^{\frac{1}{3}}(U)$.

It will be convenient to denote, for $U \subset \Omega$,

$$N_t(f, U) := \sup_{|z| \leq t} t^{-\frac{1}{3}} \|D^z f\|_{L^3(U)}$$

with this quantity we can also define the distinguished subspace, lying between $B_{3,q}^{\frac{1}{3}}$ with $q < \infty$ and $B_{3,\infty}^{\frac{1}{3}}$:

$$B_{3,c_0}^{\frac{1}{3}}(\Omega) = B_{3,\infty}^{\frac{1}{3}}(\Omega) \cap \{f : \lim_{t \rightarrow 0} N_t(f, \Omega) = 0\}.$$

In order to detect and describe the singularities of solutions of equation (M), it is customary to test the equation on suitable renormalization of the solution (see [20, 25, 15]):

Definition 2.2. A function $\Phi \in C^2(S^1, \mathbb{R}^2)$ is an entropy for the equation (M) if

$$e^{it} \cdot \frac{d}{dt} [\Phi(e^{it})] = 0 \quad \forall t \in \mathbb{R}.$$

The set of all entropies is denoted by ENT .

This definition is designed so that any smooth unit field m solving (M) satisfies $\nabla \cdot \Phi(m) = 0$ for $\Phi \in ENT$. In contrast, if m has only bounded variation, $\nabla \cdot \Phi(m)$ will be a measure concentrated on the jump set of m , called the entropy production associated to Φ . In other words, BV -type jump discontinuities of m are detected by such divergences: already in [29], in the context of the Aviles-Giga functional (AG), a special family of ‘‘cubic’’ entropies $\Sigma_{\alpha_1, \alpha_2}$ were introduced, depending on a chosen orthonormal frame of coordinates (α_1, α_2) :

$$\Sigma_{\alpha_1, \alpha_2}(z) = \frac{4}{3} \left((z \cdot \alpha_2)^3 \alpha_1 + (z \cdot \alpha_1)^3 \alpha_2 \right). \quad (2.1)$$

The maps $\Sigma_{\alpha_1, \alpha_2}$ are easily seen to belong to ENT . The divergences $\nabla \cdot \Sigma_{\alpha_1, \alpha_2}(m) = \nabla \cdot \Sigma_{\alpha_1, \alpha_2}(\nabla^\perp u)$ of these entropies detect the jump discontinuities of ∇u , according to the relative orientation of the discontinuity set $J_{\nabla u}$ with respect to the chosen frame (α_1, α_2) . An optimization procedure over the frame bundle provides the lower bound

$$\Gamma - \liminf_{\varepsilon} AG_\varepsilon(u) \geq \frac{1}{3} \int_{J_{\nabla u}} |\nabla^+ u - \nabla^- u|^3 d\mathcal{H}^1, \quad (2.2)$$

(in the $W^{1,3}$ topology) at functions u such that $\nabla u \in BV(\Omega)$ and $|\nabla u| = 1$ almost everywhere. Here $J_{\nabla u}$ is the jump set of the gradient and $\nabla^\pm u$ are its traces on $J_{\nabla u}$, see [2]. The cubic power of the jump appearing in (2.2) hints at the Besov scale $B_{3,q}^{\frac{1}{3}}$ we are considering here. For functions u with $\nabla u \in BV$, the right-hand side of (2.2) can be conveniently expressed in terms of the entropy productions, since it holds

$$\frac{1}{3} |\nabla^+ u - \nabla^- u|^3 \mathcal{H}^1 \llcorner J_{\nabla u} = \bigvee_{(\alpha_1, \alpha_2)} \|\nabla \cdot \Sigma_{\alpha_1, \alpha_2}(u)\|.$$

This is proved in [2, Theorem 3.8] (see also [25]). Here $\|\mu\|$ denotes the total variation measure of a complex-valued measure μ , and the symbol \bigvee denotes the least upper bound of a family of measures [3, Definition 1.68]:

$$\bigvee_{\alpha \in A} \mu_\alpha(E) := \sup \left\{ \sum_{\{\alpha'\} \subset A} \mu_{\alpha'}(E_{\alpha'}) : \{E_{\alpha'}\} \text{ pairwise disjoint, } E = \bigcup_{\alpha'} E_{\alpha'} \right\}.$$

Hence the estimate (2.2) provides a control of the entropy production associated to the cubic entropies (2.1) by the Aviles-Giga energy. In fact for any entropy $\Phi \in ENT$ it is shown in [20] that limits m of sequences $m_\varepsilon = \nabla^\perp u_\varepsilon$ with $AG_\varepsilon(u_\varepsilon) \leq M$ satisfy

$$\|\nabla \cdot \Phi(m)\|(\Omega) \leq C \|D^2 \Phi\|_\infty \liminf_{\varepsilon \rightarrow 0} AG_\varepsilon(u_\varepsilon).$$

In particular all the entropy productions are finite measures. This motivates the following

Definition 2.3. We say that a vector field m solving (M) has *locally finite weak entropy production* in Ω if for every $\Phi \in ENT$ we have

$$\nabla \cdot \Phi(m) \in \mathcal{M}_{\text{loc}}(\Omega). \quad (\text{wFEP})$$

If furthermore

$$\bigvee_{\Phi \in ENT, \|D^2 \Phi\|_\infty \leq 1} \|\nabla \cdot \Phi(m)\| \in \mathcal{M}_{\text{loc}}(\Omega), \quad (\text{sFEP})$$

we say that m has *locally finite strong entropy production* in Ω .

Remark 2.4. Limits of sequences $m_\varepsilon = \nabla^\perp u_\varepsilon$ with $AG_\varepsilon(u_\varepsilon) \leq M$ satisfy (sFEP).

Definition 2.5. We say that a vector field m solving (M) satisfies the kinetic formulation if there exists a Radon measure $\sigma \in \mathcal{M}_{\text{loc}}(\Omega \times \mathbb{R}/2\pi\mathbb{Z})$ such that

$$e^{is} \cdot \nabla_x \mathbb{1}_{e^{is} \cdot m(x) > 0} = \partial_s \sigma \quad \text{in } \mathcal{D}'(\Omega \times \mathbb{R}/2\pi\mathbb{Z}). \quad (\text{KIN})$$

The main Theorem of this paper is the following:

Theorem 2.6. *Let m satisfy (M). The following four conditions are equivalent:*

- (i) m has locally finite weak entropy production, (wFEP);
- (ii) m satisfies the kinetic equation (KIN);
- (iii) $m \in B_{3,\infty,\text{loc}}^{1/3}(\Omega)$;
- (iv) m has locally finite strong entropy production, (sFEP).

Remark 2.7. It is interesting to recall the following boundary behaviour of solutions of conservation laws with finite entropy production [41, Theorem 1.1] (see also [13, Theorem 2.5]): the field m admit a strong L^1 trace on $\partial\Omega$, in the sense that there exists a function $v \in L^\infty(\partial\Omega, S^1)$ such that

$$\text{ess lim}_{s \rightarrow 0} \int_{\partial\Omega} |u(\psi(s, x)) - v(x)| d\mathcal{H}^1(x) = 0,$$

where ψ is a suitable parametrization of a neighborhood of $\partial\Omega$.

3. PROOF OF THE MAIN THEOREM

The proof of Theorem 2.6 is divided into three propositions. The implication (iv) \Rightarrow (i) is trivial.

3.1. Finite entropy implies kinetic formulation.

Proposition 3.1. *If m has weak finite entropy production (wFEP), then it satisfies the kinetic formulation (KIN).*

We will need to construct a suitable family of entropies Φ_f parametrized (linearly) by continuous functions on S^1 :

$$C^0(\mathbb{R}/2\pi\mathbb{Z}, \mathbb{R}) \ni f \mapsto \Phi_f \in ENT.$$

The construction is done in several steps. First define $\tilde{f} \in C^0(\mathbb{R}/2\pi\mathbb{Z}, \mathbb{R})$ by removing the null and the first Fourier modes:

$$\tilde{f}(t) = f(t) - \left(\frac{1}{2\pi} \int_0^{2\pi} f(s) ds \right) - \left(\frac{1}{\pi} \int_0^{2\pi} f(s) \cos s ds \right) \cos t - \left(\frac{1}{\pi} \int_0^{2\pi} f(s) \sin s ds \right) \sin t.$$

Then define $\psi_f \in C^1(\mathbb{R}/2\pi\mathbb{Z}, \mathbb{R})$ by

$$\psi_f(t) = \int_0^t \tilde{f}(s) ds.$$

Note that ψ_f is 2π -periodic since $\int_0^{2\pi} \tilde{f} = 0$. Moreover it holds

$$\int_0^{2\pi} \psi_f(s) e^{is} ds = 0,$$

since $\int_0^{2\pi} \tilde{f}(s) e^{is} ds = 0$. This allows us to define $\varphi_f \in C^2(\mathbb{R}/2\pi\mathbb{Z}, \mathbb{R}^2)$ by

$$\varphi_f(t) = \int_0^t \psi_f(s) i e^{is} ds.$$

Finally define $\Phi_f \in C^2(S^1, \mathbb{R}^2)$ by

$$\Phi_f(e^{it}) = -i\varphi_f(t - \pi/2) + i\varphi_f(t + \pi/2).$$

Then it holds

$$\begin{aligned} e^{it} \cdot \frac{d}{dt} [\Phi_f(e^{it})] &= e^{it} \cdot \left(-i\psi_f(t - \pi/2) i e^{i(t-\pi/2)} + i\psi_f(t + \pi/2) i e^{i(t+\pi/2)} \right) \\ &= -(\psi_f(t - \pi/2) + \psi_f(t + \pi/2)) e^{it} \cdot (i e^{it}) = 0, \end{aligned}$$

so that $\Phi_f \in ENT$. Note that the map $f \mapsto \Phi_f$ is linear, and that $\|\Phi_f\|_{C^2} \leq C\|f\|_{C^0}$ for some constant $C > 0$.

Remark 3.2. Note that $\Phi_{\cos(2t)} = -\frac{1}{2}\Sigma_{e_1, e_2}$ and that $\Phi_{\sin(2t)} = -\frac{1}{2}\Sigma_{\varepsilon_1, \varepsilon_2}$, where (e_1, e_2) is the standard basis and $(\varepsilon_1, \varepsilon_2)$ is its rotation by $\pi/4$. In particular, the classical entropies for the Aviles-Giga functional discovered by Jin and Kohn are parametrized by the first nontrivial modes of f (those with wavenumber 2).

The reason for defining the family of entropies $\{\Phi_f\}$ as above lies in its connection to the left-hand side of the kinetic formulation:

Lemma 3.3. *Let $\nu := e^{it} \cdot \nabla_x (\mathbf{1}_{e^{it} \cdot m(x) > 0}) \in \mathcal{D}'(\Omega \times \mathbb{R}/2\pi\mathbb{Z})$. For any $f \in C^\infty(\mathbb{R}/2\pi\mathbb{Z}, \mathbb{R})$ and $\zeta \in C_c^\infty(\Omega)$ it holds*

$$\langle \nu, \zeta \otimes \psi_f \rangle = -\langle \nabla \cdot \Phi_f(m), \zeta \rangle.$$

Proof. We have

$$\langle \nu, \psi_f(t) \zeta(x) \rangle = - \int_{\Omega} \nabla \zeta(x) \cdot \int_{\mathbb{R}/2\pi\mathbb{Z}} \psi_f(t) \mathbf{1}_{e^{it} \cdot m > 0} e^{it} dt dx,$$

and for all $x \in \Omega$, writing $m(x) = e^{i\alpha}$ we compute

$$\begin{aligned} \int_{\mathbb{R}/2\pi\mathbb{Z}} \psi_f(t) \mathbb{1}_{e^{it} \cdot m > 0} e^{it} dt &= \int_{\alpha-\pi/2}^{\alpha+\pi/2} \psi_f(t) e^{it} dt \\ &= \int_{\alpha-\pi/2}^{\alpha+\pi/2} \frac{1}{i} \varphi'_f(t) dt \\ &= -i\varphi_f(\alpha + \pi/2) + i\varphi_f(\alpha - \pi/2) \\ &= \Phi_f(m), \end{aligned}$$

hence $\langle \nu, \zeta \otimes \psi_f \rangle = - \int_{\Omega} \Phi_f(m) \cdot \nabla \zeta dx$. \square

The next lemma provides the measure σ appearing in the right-hand side of the kinetic formulation: as in [34, Theorem 3.1.6], the entropy production of the solution u of a conservation law under a certain entropy S can be written as an integral of S'' against the so-called entropy measure. In our case, observe that Φ_f is obtained by integrating f twice.

Lemma 3.4. *If m has locally finite weak entropy production in Ω , then there exists $\sigma \in \mathcal{M}_{\text{loc}}(\Omega \times \mathbb{R}/2\pi\mathbb{Z})$ satisfying*

$$\langle \nabla \cdot \Phi_f(m), \zeta \rangle = \iint_{\Omega \times \mathbb{R}/2\pi\mathbb{Z}} f(t) \zeta(x) d\sigma(x, t), \quad (3.1)$$

for every $\zeta \in C_c^\infty(\Omega)$ and every $f \in C^0(\mathbb{R}/2\pi\mathbb{Z}, \mathbb{R})$.

Proof. We consider, for any fixed $\zeta \in C_c^\infty(\Omega)$, the linear functional $T_\zeta: C^0(\mathbb{R}/2\pi\mathbb{Z}, \mathbb{R}) \rightarrow \mathbb{R}$ given by

$$T_\zeta(f) = \langle \nabla \cdot \Phi_f(m), \zeta \rangle.$$

Each functional T_ζ is continuous, since

$$|T_\zeta(f)| \leq \|\Phi_f\|_\infty \|\nabla \zeta\|_\infty \leq C \|\nabla \zeta\|_\infty \|f\|_\infty.$$

On the other hand, for any $U \subset\subset \Omega$ and $f \in C^0(\mathbb{R}/2\pi\mathbb{Z}, \mathbb{R})$, by (wFEP) it holds

$$|T_\zeta(f)| \leq \|\nabla \cdot \Phi_f(m)\|_{\mathcal{M}(U)} \quad \forall \zeta \in C_c^\infty(U), \|\zeta\|_\infty \leq 1.$$

Applying Banach-Steinhaus' theorem we deduce the existence of $C(U) > 0$ such that

$$|\langle \nabla \cdot \Phi_f(m), \zeta \rangle| \leq C(U) \|f\|_\infty \|\zeta\|_\infty,$$

for all $f \in C^0(\mathbb{R}/2\pi\mathbb{Z}, \mathbb{R})$ and $\zeta \in C_c^\infty(U)$. Since tensor products are dense in $C_c^0(\Omega \times \mathbb{R}/2\pi\mathbb{Z})$, by Riesz' representation theorem this implies the existence of $\sigma \in \mathcal{M}_{\text{loc}}(\Omega \times \mathbb{R}/2\pi\mathbb{Z})$ satisfying (3.1). \square

By Lemma 3.3 and 3.4 above, and since by definition $f = \psi'_f$, we have

$$\langle \nu - \partial_t \sigma, \psi_f(t) \zeta(x) \rangle = 0 \quad \forall f \in C^\infty(\mathbb{R}/2\pi\mathbb{Z}, \mathbb{R}), \zeta \in C_c^\infty(\Omega).$$

However ψ_f cannot be any arbitrary function $\psi \in C^\infty(\mathbb{R}/2\pi\mathbb{Z}, \mathbb{R})$. In fact it holds

$$\{\psi_f: f \in C^\infty(\mathbb{R}/2\pi\mathbb{Z})\} = \left\{ \psi \in C^\infty(\mathbb{R}/2\pi\mathbb{Z}): \psi(0) = 0 \text{ and } \int_0^{2\pi} \psi_f(s) e^{is} ds = 0 \right\}.$$

In other words, we have thus far determined ν up to the Fourier modes $\{1, \cos t, \sin t\}$ in the t -variable. The next lemma takes care of those modes.

Lemma 3.5. *For all $\zeta \in C_c^\infty(\Omega)$ it holds*

$$\langle \nu, \zeta(x) \rangle = \langle \nu, \zeta(x) \cos t \rangle = \langle \nu, \zeta(x) \sin t \rangle = 0.$$

Proof. We compute

$$\begin{aligned}
\langle \nu, \zeta(x) \rangle &= - \int_{\Omega} \nabla \zeta(x) \cdot \int_{\mathbb{R}/2\pi\mathbb{Z}} \mathbb{1}_{e^{it} \cdot m > 0} e^{it} dt dx \\
&= - \int_{\Omega} \nabla \zeta(x) \cdot (2m(x)) dx \\
&= 2 \langle \nabla \cdot m, \zeta \rangle = 0, \\
\langle \nu, \zeta(x) \cos t \rangle &= - \int_{\Omega} \nabla \zeta(x) \cdot \int_{\mathbb{R}/2\pi\mathbb{Z}} \cos t \mathbb{1}_{e^{it} \cdot m > 0} e^{it} dt dx \\
&= \int_{\Omega} \nabla \zeta(x) \cdot \left(\frac{\pi}{2} \right) dx = 0,
\end{aligned}$$

and similarly $\langle \nu, \zeta(x) \sin t \rangle = 0$. \square

Remark 3.6. A similar computation shows that

$$\langle \nu, \zeta(x) \cos((2k+1)t) \rangle = \langle \nu, \zeta(x) \sin((2k+1)t) \rangle = 0 \quad \forall k \in \mathbb{N}.$$

Therefore the measure σ does not have odd frequency Fourier modes. It can also be checked directly that for $f(t) = \cos((2k+1)t)$ and $f(t) = \sin((2k+1)t)$ it holds $\Phi_f \equiv 0$, which implies the same conclusion.

Proof of Proposition 3.1. For $f(t) = \cos t$ or $f(t) = \sin t$ we have $\tilde{f} = 0$ and therefore $\Phi_f = 0$. By Lemma 3.4 this implies

$$\langle \partial_t \sigma, \psi(t) \zeta(x) \rangle = 0 \quad \text{for } \psi(t) = 1 \text{ or } \cos t \text{ or } \sin t.$$

We deduce that

$$\langle \nu - \partial_t \sigma, \psi(t) \zeta(x) \rangle = 0,$$

for any $\zeta \in C_c^\infty(\Omega)$ and $\psi \in C^\infty(\mathbb{R}/2\pi\mathbb{Z})$, which proves (KIN). \square

3.2. Kinetic formulation implies Besov regularity.

Proposition 3.7. *If m satisfies the kinetic equation (KIN), then it belongs to $B_{3,\infty;\text{loc}}^{1/3}(\Omega)$.*

The proof of Theorem 2.6 is inspired from the kinetic averaging lemma in [23] for 1D scalar conservation laws, and the way it is revisited in [21]. Following [21, 14] we make use of the following quantity to control spatial increments of m at a fixed scale h in the direction e . Let $m : \Omega \rightarrow S^1$ be measurable: given $h > 0$, $|e| = 1$ and $x \in \Omega$ we set

$$\Delta(x, h, e) = \iint_{S^1 \times S^1} \varphi(\xi, \eta) (\xi \wedge \eta) D_e^h \chi(x, \xi) D_e^h \chi(x, \eta) d\xi d\eta,$$

where $\chi(x, \xi) = \mathbb{1}_{\xi \cdot m(x) > 0}$, $D_e^h \chi(x, \cdot) = D^{he} \chi(x, \cdot) = \chi(x + he, \cdot) - \chi(x, \cdot)$ and

$$\varphi(\xi, \eta) = (\mathbb{1}_{\xi \cdot \eta > 0} - \mathbb{1}_{\xi \cdot \eta < 0})(\mathbb{1}_{\xi \wedge \eta > 0} - \mathbb{1}_{\xi \wedge \eta < 0}). \quad (3.2)$$

The next Lemma describes the coerciveness properties of the function $\Delta(x, h, e)$, with respect to the averaged quantities

$$\frac{1}{2} \int_{S^1} \xi \chi(x + he, \xi) d\xi = m(x + he) \quad \text{and} \quad \frac{1}{2} \int_{S^1} \xi \chi(x, \xi) d\xi = m(x).$$

Lemma 3.8. *Given $m : \Omega \rightarrow S^1$, $x \in \Omega$ and $0 < h < \text{dist}(x, \partial\Omega)$, it holds:*

$$\Delta(x, h, e) \gtrsim |m(x + he) - m(x)|^3 = |D_e^h m(x)|^3.$$

Proof. It holds $\Delta(x, h, e) = \Xi(m(x + he), m(x))$ where

$$\Xi(m_1, m_2) = \iint_{S^1 \times S^1} \varphi(\xi, \eta)(\xi \wedge \eta) (\mathbf{1}_{\xi \cdot m_1 > 0} - \mathbf{1}_{\xi \cdot m_2 > 0}) (\mathbf{1}_{\eta \cdot m_1 > 0} - \mathbf{1}_{\eta \cdot m_2 > 0}).$$

Therefore it suffices to prove that

$$\Xi(m_1, m_2) \gtrsim |m_1 - m_2|^3 \quad \forall m_1, m_2 \in S^1.$$

It is easily checked that $\Xi(m_1, m_2) = \Xi(m_2, m_1)$ and, since

$$\varphi(R\xi, R\eta) = \varphi(\xi, \eta) \quad \forall R \in SO(2), \quad (3.3)$$

that $\Xi(Rm_1, Rm_2) = \Xi(m_1, m_2)$ for all $R \in SO(2)$. Therefore it is enough to consider the case $m_1 = e^{-i\beta}$, $m_2 = e^{i\beta}$ for some $\beta \in [0, \pi/2]$ and to prove

$$\Xi(e^{-i\beta}, e^{i\beta}) \gtrsim \beta^3 \quad \forall \beta \in [0, \pi/2].$$

The function φ defined in (3.2) that appears in the definition of Ξ satisfies

$$\varphi(e^{i\theta}, e^{i\psi}) = \tilde{\varphi}(\psi - \theta), \quad \tilde{\varphi}(\omega) = \mathbf{1}_{\omega \in (0, \pi/2) \bmod \pi} - \mathbf{1}_{\omega \in (\pi/2, \pi) \bmod \pi}.$$

We compute

$$\begin{aligned} \Xi(e^{-i\beta}, e^{i\beta}) &= \int_{-\pi}^{\pi} \int_{-\pi}^{\pi} \tilde{\varphi}(\psi - \theta) \sin(\psi - \theta) \\ &\quad \cdot (\mathbf{1}_{e^{i\theta} \cdot e^{-i\beta} > 0} - \mathbf{1}_{e^{i\theta} \cdot e^{i\beta} > 0}) (\mathbf{1}_{e^{i\psi} \cdot e^{-i\beta} > 0} - \mathbf{1}_{e^{i\psi} \cdot e^{i\beta} > 0}) d\theta d\psi \\ &= \int_{-\pi}^{\pi} \tilde{\varphi}(\omega) \sin(\omega) \gamma(\omega) d\omega, \end{aligned}$$

$$\text{where } \gamma(\omega) = \int_{-\pi}^{\pi} \bar{\chi}(\theta) \bar{\chi}(\theta + \omega) d\theta,$$

$$\text{and } \bar{\chi}(\theta) = \mathbf{1}_{e^{i\theta} \cdot e^{-i\beta} > 0} - \mathbf{1}_{e^{i\theta} \cdot e^{i\beta} > 0}.$$

Note that $\bar{\chi}(\theta + \pi) = \bar{\chi}(-\theta) = -\bar{\chi}(\theta)$ for almost every $\theta \in \mathbb{R}$. Therefore $\omega \mapsto \tilde{\varphi}(\omega) \sin(\omega) \gamma(\omega)$ is π -periodic and even, and

$$\Xi(e^{-i\beta}, e^{i\beta}) = 4 \int_0^{\pi/2} \tilde{\varphi}(\omega) \sin(\omega) \gamma(\omega) d\omega. \quad (3.4)$$

Moreover the integrand defining γ is π -periodic in θ , hence for all $\omega \in (0, \pi/2)$ we have

$$\gamma(\omega) = 2 \int_0^{\pi} \bar{\chi}(\theta) \bar{\chi}(\theta + \omega) d\theta.$$

Assume first $\beta \in [0, \pi/4]$. Then for $\theta \in (0, \pi)$ it holds

$$\bar{\chi}(\theta) \bar{\chi}(\theta + \omega) = \begin{cases} \mathbf{1}_{\theta \in [\pi/2 - \beta, \pi/2 + \beta - \omega]} & \text{if } \omega \in [0, 2\beta], \\ 0 & \text{if } \omega \in [2\beta, \pi/2], \end{cases}$$

and we find

$$\begin{aligned} \gamma(\omega) &= \begin{cases} 2 \cdot (2\beta - \omega) & \text{if } \omega \in [0, 2\beta], \\ 0 & \text{if } \omega \in [2\beta, \pi/2]. \end{cases} \\ &= 2 \cdot (2\beta - \omega)_+ \quad \forall \omega \in [0, \pi/2], \end{aligned}$$

Plugging this into (3.4) we deduce

$$\begin{aligned} \Xi(e^{-i\beta}, e^{i\beta}) &= 8 \int_0^{2\beta} (2\beta - \omega) \sin \omega d\omega \\ &= 8 \cdot (2\beta - \sin(2\beta)) \gtrsim \beta^3, \end{aligned}$$

for all $\beta \in [0, \pi/4]$.

Consider now $\beta \in [\pi/4, \pi/2]$. For $\theta \in [0, \pi]$ we have

$$\bar{\chi}(\theta)\bar{\chi}(\theta + \omega) = \begin{cases} \mathbb{1}_{\theta \in [\pi/2 - \beta, \pi/2 + \beta - \omega]} & \text{if } \omega \in [0, \pi - 2\beta] \\ \mathbb{1}_{\theta \in [\pi/2 - \beta, \pi/2 + \beta - \omega]} - \mathbb{1}_{\theta \in (3\pi/2 - \beta - \omega, \pi/2 + \beta)} & \text{if } \omega \in [\pi - 2\beta, \pi/2], \end{cases}$$

and therefore

$$\begin{aligned} \gamma(\omega) &= \begin{cases} 2 \cdot (2\beta - \omega) & \text{if } \omega \in [0, \pi - 2\beta], \\ 2 \cdot (\pi - 2\omega) & \text{if } \omega \in [\pi - 2\beta, \pi/2], \end{cases} \\ \Xi(e^{-i\beta}, e^{i\beta}) &= 8 \int_0^{\pi - 2\beta} (2\beta - \omega) \sin \omega \, d\omega + 8 \int_{\pi - 2\beta}^{\pi/2} (\pi - 2\omega) \sin \omega \, d\omega \\ &= 8 \cdot (2\beta - (\pi - 4\beta) \cos(2\beta) - \sin(2\beta)) \\ &\quad + 8 \cdot (2 \sin(2\beta) - 2 + (\pi - 4\beta) \cos(2\beta)) \\ &= 8 \cdot (\sin(2\beta) + 2\beta - 2) \geq 8 \left(\frac{\pi}{2} - 1 \right) \gtrsim \beta^3, \end{aligned}$$

for all $\beta \in [\pi/4, \pi/2]$. \square

To obtain bounds for the integral of $\Delta(x, h, e)$ on Ω , when m satisfies the kinetic formulation (KIN), we use the following lemma, that estimates its derivative with respect to h :

Lemma 3.9. *Suppose that m satisfies the kinetic formulation of the eikonal equation (KIN), and that $\Omega' \subset\subset \Omega'' \subset\subset \Omega$. We then have for all unit vectors e and $|h| \lesssim \text{dist}(\Omega', \partial\Omega'')$:*

$$\int_{\Omega'} \Delta(x, h, e) dx \lesssim |h| (1 + \|\sigma\|_{\mathcal{M}(\Omega'' \times \mathbb{R}/2\pi\mathbb{Z})}), \quad (3.5)$$

where the multiplicative constant depends on the distance between Ω' and $\partial\Omega''$.

Proof. Assume (KIN):

$$e^{is} \cdot \nabla_x \mathbb{1}_{e^{is} \cdot m(x) > 0} = \partial_s \sigma \quad \text{in } \mathcal{D}'(\Omega \times \mathbb{R}/2\pi\mathbb{Z}),$$

and let us assume to have intermediate domains $\tilde{\Omega}, \Omega''$ with $\Omega' \subset\subset \tilde{\Omega} \subset\subset \Omega'' \subset\subset \Omega$ and such that the distances among the boundaries of the first three are comparable. We perform the calculation of $\partial_h \int_{\Omega'} \Delta(x, h, e) dx$ for a regularized integrand, namely:

- we regularize the equation (KIN) by convolving with respect to x with a smooth approximation of the identity ρ_ε :

$$e^{is} \cdot \nabla_x \chi_\varepsilon(x, e^{is}) = \partial_s \sigma_\varepsilon, \quad \chi_\varepsilon = (\mathbb{1}_{e^{is} \cdot m(x) > 0}) * \rho_\varepsilon, \quad \sigma_\varepsilon = \sigma * \rho_\varepsilon; \quad (3.6)$$

here $\varepsilon < \text{dist}(\Omega, \partial\Omega'')$;

- we approximate φ (3.2) by a smooth φ_δ . The calculations below are valid for a generic φ and only use the skew-symmetry property $\varphi(\xi, \eta) = -\varphi(\eta, \xi)$. Assuming in addition the $SO(2)$ invariance property (3.3), and parametrizing with the angle between ξ and η , these conditions amount to require that $\tilde{\varphi}: s \mapsto \varphi(1, e^{is})$ is odd and 2π periodic. In turn, a convolution on the real line with a smooth even kernel, at scale δ , preserves both these properties. Explicitly, we set

$$\varphi_\delta(e^{i\theta}, e^{i\psi}) = \tilde{\varphi}_\delta(\psi - \theta), \quad \tilde{\varphi}_\delta = \tilde{\varphi} * \rho_\delta,$$

for some smooth even kernel ρ . This approximation has the following properties:

$$\tilde{\varphi}_\delta \rightarrow \tilde{\varphi} \text{ a.e.}, \quad |\tilde{\varphi}_\delta| \leq 1, \quad \|\tilde{\varphi}'_\delta\|_{L^1(\mathbb{R}/2\pi\mathbb{Z})} \leq 8. \quad (3.7)$$

The explicit dependence of the function Δ on the parameters ε, δ is omitted in the first calculations.

We assume without loss of generality that $e = e_1$ and use the notations $\chi^h(x, \xi) = \chi(x + he_1, \xi)$ and $D_1^h \chi = \chi^h - \chi$. Let $x \in \tilde{\Omega}$ and $|h| < \text{dist}(\tilde{\Omega}, \partial\Omega'')$. Using the skew-symmetry of φ , we have

$$\begin{aligned} \frac{\partial}{\partial h} \Delta(x, e, h) &= \frac{\partial}{\partial h} \iint_{S^1 \times S^1} \varphi(\xi, \eta) (\xi \wedge \eta) D_1^h \chi(x, \xi) D_1^h \chi(x, \eta) d\xi d\eta \\ &= \iint_{S^1 \times S^1} \varphi(\xi, \eta) (\xi \wedge \eta) [\partial_1 \chi^h(x, \xi) D_1^h \chi(x, \eta) + \partial_1 \chi^h(x, \eta) D_1^h \chi(x, \xi)] d\xi d\eta \\ &= 2 \iint_{S^1 \times S^1} \varphi(\xi, \eta) (\xi_1 \eta_2) [\partial_1 \chi^h(x, \xi) D_1^h \chi(x, \eta) + \partial_1 \chi^h(x, \eta) D_1^h \chi(x, \xi)] d\xi d\eta. \end{aligned}$$

Letting $\nu(x, e^{is}) := \partial_s \sigma(x, s)$, we use the the equation (KIN) in the form

$$\xi_1 \partial_1 \chi(x, \xi) + \xi_2 \partial_2 \chi(x, \xi) = \nu(x, \xi), \quad (3.8)$$

to replace $\xi_1 \partial_1 \chi^h(x, \xi)$ in the above and obtain

$$\begin{aligned} \frac{\partial}{\partial h} \Delta(x, e, h) &= 2 \iint_{S^1 \times S^1} \varphi(\xi, \eta) \eta_2 [(\nu^h(x, \xi) - \xi_2 \partial_2 \chi(x, \xi)) D_1^h \chi(x, \eta) + \xi_1 \partial_1 \chi^h(x, \eta) D_1^h \chi(x, \xi)] d\xi d\eta \\ &= 2 \iint_{S^1 \times S^1} \varphi(\xi, \eta) \eta_2 [(\nu^h(x, \xi) - \xi_2 \partial_2 \chi(x, \xi)) D_1^h \chi(x, \eta) - \xi_1 \chi^h(x, \eta) \partial_1 D_1^h \chi(x, \xi)] d\xi d\eta \\ &\quad + \partial_1 \left[2 \iint_{S^1 \times S^1} \varphi(\xi, \eta) \eta_2 \xi_1 \chi^h(x, \eta) D_1^h \chi(x, \xi) d\xi d\eta \right] =: I_1 + \partial_1 A_1. \end{aligned}$$

The term $\partial_1 A_1$ is a boundary term and will be treated at the end. Focusing on I_1 , we can expand $\xi_1 \partial_1 D_1^h \chi(x, \xi)$ and use (3.8) to deduce

$$\begin{aligned} I_1 &= 2 \iint_{S^1 \times S^1} \varphi(\xi, \eta) \eta_2 [(\nu^h(x, \xi) - \xi_2 \partial_2 \chi(x, \xi)) (\chi^h(x, \eta) - \chi(x, \eta)) \\ &\quad - (\nu^h(x, \xi) - \nu(x, \xi) - \xi_2 \partial_2 \chi^h(x, \xi) + \xi_2 \partial_2 \chi(x, \xi)) \chi^h(x, \eta)] d\xi d\eta \\ &= 2 \iint_{S^1 \times S^1} \varphi(\xi, \eta) \eta_2 [-\nu^h(x, \xi) \chi(x, \eta) + \nu(x, \xi) \chi^h(x, \eta) \\ &\quad + \xi_2 \partial_2 \chi^h(x, \xi) \chi(x, \eta) - \xi_2 \partial_2 \chi(x, \xi) \chi^h(x, \eta)] d\xi d\eta \\ &= 2 \iint_{S^1 \times S^1} \varphi(\xi, \eta) \eta_2 [-\nu^h(x, \xi) \chi(x, \eta) + \nu(x, \xi) \chi^h(x, \eta)] \\ &\quad + 2 \iint_{S^1 \times S^1} \varphi(\xi, \eta) \eta_2 \xi_2 [\partial_2 \chi^h(x, \xi) \chi(x, \eta) - \partial_2 \chi(x, \xi) \chi^h(x, \eta)] d\xi d\eta. \end{aligned}$$

Exchanging ξ and η only in the last term of the second integral, we can rewrite

$$\begin{aligned} I_1 &= 2 \iint_{S^1 \times S^1} \varphi(\xi, \eta) \eta_2 [-\nu^h(x, \xi) \chi(x, \eta) + \nu(x, \xi) \chi^h(x, \eta)] \\ &\quad + 2 \iint_{S^1 \times S^1} \varphi(\xi, \eta) \eta_2 \xi_2 [\partial_2 \chi^h(x, \xi) \chi(x, \eta) + \partial_2 \chi(x, \eta) \chi^h(x, \xi)] d\xi d\eta \\ &= 2 \iint_{S^1 \times S^1} \varphi(\xi, \eta) \eta_2 [-\nu^h(x, \xi) \chi(x, \eta) + \nu(x, \xi) \chi^h(x, \eta)] \\ &\quad + \partial_2 \left[2 \iint_{S^1 \times S^1} \varphi(\xi, \eta) \eta_2 \xi_2 \chi(x, \eta) \chi^h(x, \xi) d\xi d\eta \right] =: I_2 + \partial_2 A_2. \end{aligned}$$

Therefore we have $\frac{\partial}{\partial h} \Delta(x, h, e) = I_2 + \partial_1 A_1 + \partial_2 A_2$, where

$$|A| = |(A_1, A_2)| \leq 8\pi \quad \text{pointwise } \forall(x, h, e).$$

The most important term in the estimate is I_2 , since the extra term is a divergence $\nabla_x \cdot A$ of a bounded vectorfield, hence it can be treated as a boundary term. In polar coordinates I_2 becomes:

$$\begin{aligned} I_2 &= 2 \iint_{[0, 2\pi[\times [0, 2\pi[} \tilde{\varphi}(\psi - \theta) \sin \psi [-\partial_\theta \sigma^h(x, \theta) \chi(x, e^{i\psi}) + \partial_\theta \sigma(x, \theta) \chi^h(x, e^{i\psi})] d\theta d\psi \\ &= 2 \iint_{[0, 2\pi[\times [0, 2\pi[} \tilde{\varphi}'(\psi - \theta) \sin \psi [-\sigma^h(x, \theta) \chi(x, e^{i\psi}) + \sigma(x, \theta) \chi^h(x, e^{i\psi})] d\theta d\psi. \end{aligned}$$

Recall now that the above was derived for an approximation φ_δ of φ and for a solution of the regularized kinetic equation (3.6) at scale ε . Writing this dependence explicitly we have:

$$\begin{aligned} I_2 &= I_2^{\varepsilon, \delta} = -2 \int_0^{2\pi} \sigma_\varepsilon^h(x, e^{i\theta}) \int_0^{2\pi} \tilde{\varphi}'_\delta(\psi - \theta) \chi_\varepsilon(x, e^{i\psi}) \sin \psi d\psi d\theta \\ &\quad + 2 \int_0^{2\pi} \sigma_\varepsilon(x, e^{i\theta}) \int_0^{2\pi} \tilde{\varphi}'_\delta(\psi - \theta) \chi_\varepsilon^h(x, e^{i\psi}) \sin \psi d\psi d\theta. \end{aligned}$$

Recalling (3.7) and the fact that $|\chi_\varepsilon| \leq 1$ a.e., we deduce

$$|I_2^{\varepsilon, \delta}| \lesssim \int_0^{2\pi} (|\sigma_\varepsilon^h(x, \theta)| + |\sigma_\varepsilon(x, \theta)|) d\theta.$$

Plugging this estimate into the identity $\partial_h \Delta^{\varepsilon, \delta} = I_2^{\varepsilon, \delta} + \nabla_x \cdot A^{\varepsilon, \delta}$ yields

$$\frac{\partial}{\partial h} \Delta^{\varepsilon, \delta}(x, h, e) \lesssim \int_0^{2\pi} (|\sigma_\varepsilon^h(x, \theta)| + |\sigma_\varepsilon(x, \theta)|) d\theta + \nabla_x \cdot A^{\varepsilon, \delta}.$$

Recalling that $A^{\varepsilon, \delta}$ is a uniformly bounded vector field, we may test the above against any nonnegative $\gamma \in C_c^\infty(\tilde{\Omega})$ and obtain

$$\begin{aligned} \frac{\partial}{\partial h} \int_\Omega \gamma(x) \Delta^{\varepsilon, \delta}(x, h, e) dx &\lesssim \|\gamma\|_{C^0} \left(\|\sigma_\varepsilon^h\|_{L^1(\tilde{\Omega} \times \mathbb{R}/2\pi\mathbb{Z})} + \|\sigma_\varepsilon\|_{L^1(\tilde{\Omega} \times \mathbb{R}/2\pi\mathbb{Z})} \right) + \|\nabla \gamma\|_{C^0} \\ &\lesssim \|\gamma\|_{C^0} \|\sigma\|_{\mathcal{M}(\Omega'' \times \mathbb{R}/2\pi\mathbb{Z})} + \|\nabla \gamma\|_{C^0}, \end{aligned}$$

for $|h| + \varepsilon \leq \text{dist}(\tilde{\Omega}, \partial\Omega'')$ and $\delta > 0$. Integrating with respect to h we find that

$$\frac{1}{|h|} \int_\Omega \gamma(x) \Delta^{\varepsilon, \delta}(x, h, e) dx \lesssim \|\gamma\|_{C^0} \|\sigma\|_{\mathcal{M}(\Omega'' \times \mathbb{R}/2\pi\mathbb{Z})} + \|\nabla \gamma\|_{C^0}.$$

By dominated convergence we may pass to the limit $\varepsilon, \delta \rightarrow 0$ in the left-hand side. Then it remains to choose $\gamma \equiv 1$ in Ω' to obtain the claimed estimate (3.5). \square

We can now prove Proposition 3.7:

Proof. The proof follows combining the results of Lemmas 3.8 and 3.9. For $t < \text{dist}(\Omega', \partial\Omega'')$,

$$\begin{aligned} [N_t(m, \Omega')]^3 &= \frac{1}{t} \sup_{|e|=1, |h| \leq t} \int_{\Omega'} |D_e^h m(x)|^3 dx \\ &\lesssim \frac{1}{t} \sup_{|e|=1, |h| \leq t} \int_{\Omega'} \Delta(x, e, h) dx \\ &\lesssim 1 + \|\sigma\|_{\mathcal{M}(\Omega'' \times \mathbb{R}/2\pi\mathbb{Z})}. \end{aligned}$$

For other values of t up to 1, the triangular inequality and the boundedness of m yield a trivial control on $N_t(f, \Omega)$. Together, these estimates give the desired bound on the local Besov norm $B_{3, \infty}^{1/3}(\Omega')$. \square

3.3. Besov regularity implies finite entropy. For the proofs of the next propositions and lemmas, we let $m_\varepsilon := m * \rho_\varepsilon$ be a regularization with a standard kernel (with $\text{spt}(\rho) \subset B_1$ and $\nabla \rho = 0$ in $B_{1/2}$).

Proposition 3.10. *If m solves (M) and belongs to the space $B_{3,\infty,\text{loc}}^{1/3}(\Omega)$, then m has locally finite strong entropy production (sFEP):*

$$\bigvee_{\Phi \in ENT, \|D^2\Phi\|_\infty \leq 1} \|\nabla \cdot \Phi(m)\|(A) \lesssim [m]_{B_{3,\infty}^{1/3}(A)}^3 \quad \text{for } A \subset\subset \Omega.$$

Proof. For a given $\Phi \in ENT$, we consider its extension $\tilde{\Phi} \in C_c^2(\mathbb{R}^2, \mathbb{R}^2)$ given in polar coordinates by $\tilde{\Phi}(re^{i\theta}) = \eta(r)\Phi(e^{i\theta})$, where η is a fixed cut-off function $\eta \in C_c^\infty(0, \infty)$ satisfying $\eta \equiv 0$ outside $(1/2, 2)$ and $\eta(1) = 1$.

Following [20, 15], for the mollified field m_ε we can single out in the entropy production the contribution of the radial oscillation:

$$\nabla \cdot \tilde{\Phi}(m_\varepsilon) = \Psi(m_\varepsilon) \cdot \nabla(1 - |m_\varepsilon|^2),$$

where $\Psi \in C_c^1(\mathbb{R}^2, \mathbb{R}^2)$ is a regular vectorfield. Given a test function $\phi \in C_c^\infty(\Omega)$ we can integrate by parts

$$\langle \nabla \cdot \tilde{\Phi}(m_\varepsilon), \phi \rangle = - \int_\Omega \phi(x) \Psi(m_\varepsilon(x)) \cdot \nabla(1 - |m_\varepsilon(x)|^2) dx =: A_\varepsilon[\phi] + B_\varepsilon[\phi],$$

where

$$\begin{aligned} A_\varepsilon[\phi] &= \int_\Omega \nabla \phi(x) \cdot \Psi(m_\varepsilon(x)) (1 - |m_\varepsilon(x)|^2) dx, \\ B_\varepsilon[\phi] &= \int_\Omega \phi(x) \nabla \cdot [\Psi(m_\varepsilon(x))] (1 - |m_\varepsilon(x)|^2) dx \\ &= \int_\Omega \phi(x) \text{Tr}[D\Psi(m_\varepsilon(x)) \nabla m_\varepsilon(x)] (1 - |m_\varepsilon(x)|^2) dx. \end{aligned}$$

While $A_\varepsilon[\phi] \rightarrow 0$, trivially because $|m| = 1$ almost everywhere, the second integral $B_\varepsilon[\phi]$ can be bounded by

$$B_\varepsilon[\phi] \lesssim \|\phi\|_{L^\infty} \|D\Psi\|_{L^\infty} \int_{\text{spt}(\phi)} |\nabla m_\varepsilon(x)| |1 - |m_\varepsilon(x)|^2| dx.$$

Since 3 and $\frac{3}{2}$ are dual exponents, using Lemmas 3.11 and 3.12 below, the last integral can be bounded by

$$\|1 - |m_\varepsilon|^2\|_{L^{\frac{3}{2}}(\text{spt}(\phi))} \|\nabla m_\varepsilon\|_{L^3(\text{spt}(\phi))} \lesssim N_\varepsilon(m, \text{spt}(\phi))^3.$$

Noting that $|D\Psi| \lesssim |D^2\tilde{\Phi}|$ and letting $\varepsilon \rightarrow 0$ we deduce that

$$|\langle \nabla \cdot \Phi(m), \phi \rangle| \lesssim \|\phi\|_{L^\infty} \|D^2\Phi\|_\infty \liminf_{\varepsilon \rightarrow 0} N_\varepsilon(m, \text{spt}(\phi))^3,$$

and therefore

$$\|\nabla \cdot \Phi(m)\|(U) \lesssim \|D^2\Phi\|_\infty \liminf_{\varepsilon \rightarrow 0} N_\varepsilon(m, \bar{U})^3,$$

for all $U \subset\subset \Omega$. Note that $N_\varepsilon(m, \bar{U})$ involves integrals with respect to x over the sets \bar{U} and $\bar{U} + \varepsilon y$, hence given a finite family of open and distant sets $U_1, \dots, U_k \subset\subset A \subset\subset \Omega$, and a corresponding family of entropies Φ_1, \dots, Φ_k with $\|D^2\Phi_j\|_\infty \leq 1$, if ε is small enough it holds

$$\sum_j \|\nabla \cdot \Phi_j(m)\|(U_j) \lesssim \liminf_{\varepsilon \rightarrow 0} \sum_j N_\varepsilon(m, U_j)^3 = \liminf_{\varepsilon \rightarrow 0} N_\varepsilon(m, A)^3.$$

Recalling the definitions of the least upper bound measure and of the Besov seminorm, this implies the conclusion of Proposition 3.10. \square

In the proof of Proposition 3.10, we used the two following lemmas on the growth of certain norms of the regularized field m_ε . Their proof is an adaptation to the Besov scale of corresponding statements for Sobolev functions, treated in [15].

Lemma 3.11. *If $m \in B_{3,\infty,\text{loc}}^{1/3}(\Omega)$, and $\Omega' \subset\subset \Omega$, then for every $\varepsilon \lesssim \text{dist}(\Omega', \partial\Omega)$*

$$\int_{\Omega'} |\nabla m_\varepsilon|^3 dx \lesssim \varepsilon^{-2} N_\varepsilon(m, \Omega')^3.$$

Proof. As in [15, Proof of Proposition 3, Step 6(ii)], for ε small enough we have the pointwise bound:

$$|\nabla m_\varepsilon(x)| \leq \frac{\|\nabla \rho\|_\infty}{\varepsilon^3} \int_{B_\varepsilon(0) \setminus B_{\varepsilon/2}(0)} |m(x+z) - m(x)| dz.$$

Applying Jensen inequality and integrating in Ω' one obtains

$$\begin{aligned} \int_{\Omega'} |\nabla m_\varepsilon(x)|^3 dx &\lesssim \frac{1}{\varepsilon^3} \int_{\Omega'} \int_{B_\varepsilon(0) \setminus B_{\varepsilon/2}(0)} |m(x+z) - m(x)|^3 dz dx \\ &= \frac{1}{\varepsilon^3} \int_{\Omega'} \int_{B_\varepsilon(0)} |D^z m(x)|^3 dz dx \\ &\leq \frac{1}{\varepsilon^2} N_\varepsilon(m, \Omega')^3. \end{aligned}$$

□

Lemma 3.12. *With the notations of Lemma 3.11, it holds:*

$$\int_{\Omega'} (1 - |m_\varepsilon|^2)^{3/2} dx \lesssim \varepsilon N_\varepsilon(m, \Omega')^3.$$

Proof. As in [15, Proof of Proposition 3, Step 6(i)], using that $|m| = 1$ almost everywhere we obtain the pointwise bound

$$\begin{aligned} |1 - |m_\varepsilon|^2|(x) &\lesssim \int_{B_\varepsilon(x)} \int_{B_\varepsilon(x)} \rho_\varepsilon(x-y) \rho_\varepsilon(x-z) |m(y) - m(z)|^2 dy dz \\ &\lesssim \int_{B_\varepsilon(x)} \rho_\varepsilon(x-y) |m(y) - m(x)|^2 dy = \int_{B_\varepsilon(0)} \rho_\varepsilon(z) |m(x+z) - m(x)|^2 dz \end{aligned}$$

Then by Hölder's inequality

$$\begin{aligned} \int_{\Omega'} |(1 - |m_\varepsilon(x)|^2)|^{3/2} dx &\lesssim \int_{\Omega'} \int_{B_\varepsilon(0)} \rho_\varepsilon(z) |D^z m(x)|^3 dz dx \\ &\leq \sup_{|z| \leq \varepsilon} \|D^z m\|_{L^3(\Omega')}^3 \\ &\leq \varepsilon N_\varepsilon(m, \Omega')^3. \end{aligned}$$

□

To conclude the proof of Theorem 2.6 we remark that the implication (iv) \Rightarrow (i) is trivial.

4. COROLLARIES AND FURTHER COMMENTS

4.1. Sharp differentiability for zero energy states. We observe that if $m \in B_{3,\text{co},\text{loc}}^{1/3}(\Omega)$ (in particular when $m \in B_{3,q,\text{loc}}^{1/3}(\Omega)$, $q < \infty$), then thanks to Lemmas 3.11 and 3.12 we have

$$\varepsilon^{2/3} \nabla m_\varepsilon \rightarrow 0 \quad \text{in } L_{\text{loc}}^3(\Omega) \quad \text{and} \quad \varepsilon^{-2/3} (1 - |m_\varepsilon|^2) \rightarrow 0 \quad \text{in } L_{\text{loc}}^{3/2}(\Omega).$$

Therefore the conclusion of Proposition 3.10 can be refined to

$$\|\nabla \cdot \Phi(m)\| = 0 \quad \text{for every } \Phi \in ENT.$$

That is, slightly better regularity rules out entropy production. This in turn implies much stronger regularity properties: m is locally Lipschitz outside a locally finite set of vortices, [26].

4.2. The mass of the entropy measure $\|\sigma\|(\Omega)$. In the micromagnetics model studied by Rivière-Serfaty in [37], twice the total variation of the kinetic measure provides a sharp asymptotic lower bound for the energy, [37, Theorem 1]. In this paragraph we investigate whether this property holds for our model (M), at least in the BV case. Recall [12] that for $m = \nabla^\perp u \in BV(\Omega)$ satisfying (M) it holds

$$(\Gamma - \lim AG_\varepsilon)(u) = \frac{1}{3} \int_{J_m} |m^+ - m^-|^3 d\mathcal{H}^1.$$

Hence the question we raise is whether this equals $2\|\sigma\|(\Omega)$.

For a given $m \in BV(\Omega)$ satisfying (M) we compute $\|\sigma\|$ as follows. In light of Lemma 3.4 it holds

$$\|\sigma\| = \bigvee_{|f| \leq 1} \|\nabla \cdot \Phi_f(m)\|.$$

On the other hand, Remark 3.2 and the results in [2] ensure that

$$\begin{aligned} 2\|\sigma\| &= 2 \bigvee_{|f| \leq 1} \|\nabla \cdot \Phi_f(m)\| \\ &\geq \bigvee_{(\alpha_1, \alpha_2)} \|\nabla \cdot \Sigma_{\alpha_1, \alpha_2}(m)\| = \frac{1}{3} |m^+ - m^-|^3 \mathcal{H}^1 \llcorner J_m. \end{aligned} \quad (4.1)$$

Proposition 4.1. *If Dm has a nontrivial jump part, then the inequality in (4.1) is strict.*

Proof. According to [25, Theorem 3], since the set $\{\Phi_f : |f| \leq 1\}$ is symmetric (stable under multiplication by -1) and equivariant (stable under conjugation by any rotation), it holds

$$\bigvee_{|f| \leq 1} \|\nabla \cdot \Phi_f(m)\| = c(|m^+ - m^-|) \mathcal{H}^1 \llcorner J_m. \quad (4.2)$$

for a certain cost function c . This cost function is given by

$$c(s) = \sup \{ (\Phi_f(m^+) - \Phi_f(m^-)) \cdot \nu \},$$

where the supremum is taken among :

- all possible jumps $m^\pm \in S^1$ of size $|m^+ - m^-| = s$,
- all possible normal vectors $\nu \in S^1$ with the admissibility condition $(m^+ - m^-) \cdot \nu = 0$ (due to the divergence constraint $\nabla \cdot m = 0$),
- and all possible f with $|f| \leq 1$.

Using again the symmetry and equivariance of $\{\Phi_f\}$, we can simplify this as

$$c(s) = \sup \{ (\Phi_f(e^{i\beta}) - \Phi_f(e^{-i\beta})) \cdot e_1, |f| \leq 1 \} \quad \text{for } s = 2 \sin \beta, \beta \in [0, \pi/2].$$

For angles $\beta \in [-\pi/2, \pi/2]$ (we want to apply this computation also to $-\beta$) it holds

$$\begin{aligned} e_1 \cdot \Phi_f(e^{i\beta}) &= \Re(-i\varphi_f(\beta - \pi/2) + i\varphi_f(\beta + \pi/2)) \\ &= \Re \left(\int_0^{\beta - \pi/2} \psi_f(s) e^{is} ds - \int_0^{\beta + \pi/2} \psi_f(s) e^{is} ds \right) \\ &= - \int_{\beta - \pi/2}^{\beta + \pi/2} \psi_f(s) \cos s ds = - [\psi_f \sin]_{\beta - \pi/2}^{\beta + \pi/2} + \int_{\beta - \pi/2}^{\beta + \pi/2} \tilde{f}(s) \sin s ds \\ &= - \cos \beta \left(\int_0^{\beta + \pi/2} \tilde{f} + \int_0^{\beta - \pi/2} \tilde{f} \right) + \int_{\beta - \pi/2}^{\beta + \pi/2} \tilde{f} \sin \end{aligned}$$

$$\begin{aligned}
&= -\cos \beta \left[\int_0^{\beta+\pi/2} f - \frac{1}{2\pi}(\beta + \pi/2) \int_0^{2\pi} f - \frac{1}{\pi} \left(\int_0^{\beta+\pi/2} \cos \right) \int_0^{2\pi} f \cos \right. \\
&\quad - \frac{1}{\pi} \left(\int_0^{\beta+\pi/2} \sin \right) \int_0^{2\pi} f \sin + \int_0^{\beta-\pi/2} f - \frac{1}{2\pi}(\beta - \pi/2) \int_0^{2\pi} f \\
&\quad \left. - \frac{1}{\pi} \left(\int_0^{\beta-\pi/2} \cos \right) \int_0^{2\pi} f \cos - \frac{1}{\pi} \left(\int_0^{\beta-\pi/2} \sin \right) \int_0^{2\pi} f \sin \right] \\
&\quad + \int_{\beta-\pi/2}^{\beta+\pi/2} f \sin - \frac{1}{2\pi} \left(\int_{\beta-\pi/2}^{\beta+\pi/2} \sin \right) \int_0^{2\pi} f \\
&\quad - \frac{1}{\pi} \left(\int_{\beta-\pi/2}^{\beta+\pi/2} \cos \sin \right) \int_0^{2\pi} f \cos - \frac{1}{\pi} \left(\int_{\beta-\pi/2}^{\beta+\pi/2} \sin^2 \right) \int_0^{2\pi} f \sin \\
&= -\cos \beta \left[\int_0^{\beta+\pi/2} f + \int_0^{\beta-\pi/2} f - \frac{\beta}{\pi} \int_0^{2\pi} f - \frac{2}{\pi} \int_0^{2\pi} f \sin \right] \\
&\quad + \int_{\beta-\pi/2}^{\beta+\pi/2} f \sin - \frac{1}{\pi} \sin \beta \int_0^{2\pi} f - \frac{1}{2} \int_0^{2\pi} f \sin.
\end{aligned}$$

Hence for any $\beta \in [0, \pi/2]$,

$$\begin{aligned}
e_1 \cdot (\Phi_f(e^{i\beta}) - \Phi_f(e^{-i\beta})) &= \cos \beta \left[\int_0^{-\beta+\pi/2} f + \int_0^{-\beta-\pi/2} f - \int_0^{\beta+\pi/2} f - \int_0^{\beta-\pi/2} f + \frac{2\beta}{\pi} \int_0^{2\pi} f \right] \\
&\quad + \int_{\beta-\pi/2}^{\beta+\pi/2} f \sin - \int_{-\beta-\pi/2}^{-\beta+\pi/2} f \sin - \frac{2}{\pi} \sin \beta \int_0^{2\pi} f \\
&= -\cos \beta \int_{-\beta-\pi/2}^{\beta-\pi/2} f - \cos \beta \int_{-\beta+\pi/2}^{\beta+\pi/2} f + \int_{-\beta+\pi/2}^{\beta+\pi/2} f \sin - \int_{-\beta-\pi/2}^{\beta-\pi/2} f \sin \\
&\quad - \frac{2}{\pi} (\sin \beta - \beta \cos \beta) \int_0^{2\pi} f \\
&= \int_0^{2\pi} g_\beta f,
\end{aligned}$$

where g_β is π -periodic and

$$g_\beta(t) = (\sin t - \cos \beta) \mathbf{1}_{\pi/2 - \beta \leq t \leq \pi/2 + \beta} - \frac{2}{\pi} (\sin \beta - \beta \cos \beta) \quad \forall t \in [0, \pi].$$

The above computation with $f(t) = \cos(2t)$ yields an entropy production equal to $(2 \sin \beta)^3/6$, as expected. On the other hand the supremum of the above quantity over $|f| \leq 1$ is given by $\|g_\beta\|_{L^1(0,2\pi)}$. This supremum is not attained by a continuous function when $\beta > 0$. In other words, for any jump of size $s > 0$ we have $c(s) > s^3/6$. In view of (4.2) this shows that equality in (4.1) can not happen unless Dm has a trivial jump part. \square

To calculate the value of $c(s)$ we observe the following. Since g_β is π -periodic and even it holds

$$\|g_\beta\|_{L^1(0,2\pi)} = 4 \int_0^{\pi/2} |g_\beta|.$$

The function g_β is negative in $[0, t_\beta]$ and positive in $(t_\beta, \pi/2]$, where $t_\beta \in [\pi/2 - \beta, \pi/2]$ is characterized by

$$\sin t_\beta - \cos \beta = \frac{2}{\pi} (\sin \beta - \beta \cos \beta).$$

Moreover it holds that $\int_0^{\pi/2} g_\beta = 0$, hence we find

$$\begin{aligned} \int_0^{\pi/2} |g_\beta| &= \int_0^{t_\beta} (-g_\beta) + \int_{t_\beta}^{\pi/2} g_\beta = 2 \int_{t_\beta}^{\pi/2} g_\beta \\ &= 2 \cos t_\beta - 2(\pi/2 - t_\beta) \left(\cos \beta + \frac{2}{\pi} (\sin \beta - \beta \cos \beta) \right). \end{aligned}$$

With this expression it can be checked that

$$\|g_\beta\|_{L^1(0,2\pi)} \sim \frac{1}{6}(2\beta)^3 \quad \text{as } \beta \rightarrow 0,$$

hence $c(s) \sim s^3/6$ for $s \rightarrow 0$, so that the measure $\|\sigma\|$ does behave like the right-hand side of (4.1) for very small jumps.

4.3. Partial regularity obtained by using only the Jin-Kohn entropies. In this paragraph, we show how to obtain fractional differentiability of a solution m of (M) having finite entropy production for every entropy (2.1) in the class of Jin-Kohn:

$$\Sigma_{\alpha_1, \alpha_2}(z) = \frac{4}{3} \left((z \cdot \alpha_2)^3 \alpha_1 + (z \cdot \alpha_1)^3 \alpha_2 \right).$$

Recall that (α_1, α_2) is a positive orthonormal frame $(R_\theta e_1, R_\theta e_2)$, and notice moreover that every entropy is a linear combination of two basic entropies Σ_{e_1, e_2} and $\Sigma_{\varepsilon_1, \varepsilon_2}$:

$$\Sigma_{R_\theta e_1, R_\theta e_2}(z) = \cos(2\theta) \Sigma_{e_1, e_2}(z) + \sin(2\theta) \Sigma_{\varepsilon_1, \varepsilon_2}(z),$$

where $\varepsilon_1 = \frac{e_1 + e_2}{\sqrt{2}}$, $\varepsilon_2 = \frac{-e_1 + e_2}{\sqrt{2}}$.

In [30] the authors show that whenever the entropy production associated to Σ_{e_1, e_2} and $\Sigma_{\varepsilon_1, \varepsilon_2}$ vanish (which is equivalent to all the Jin-Kohn entropy productions vanishing), then in fact all entropy productions vanish and the rigidity result of [26] applies. Hence it is natural to wonder whether in general, controlling the total variation of these two basic entropy productions is enough to obtain the $B_{3, \infty}^{1/3}$ estimate (which we obtained here using all entropy productions). We do not provide an answer to this question, but show how a method described in [22] can be combined with estimates derived in [30] to obtain a $B_{4, \infty}^s$ estimate for all $s < 1/4$.

To this end we set

$$\begin{aligned} \Delta_{JK}(x, h, e) &= D_e^h \Sigma_{e_1, e_2}(m(x)) \wedge D_e^h \Sigma_{\varepsilon_1, \varepsilon_2}(m(x)) \\ &= \det D_e^h(\Sigma_{e_1, e_2}(m), \Sigma_{\varepsilon_1, \varepsilon_2}(m))(x). \end{aligned}$$

Here we recall that D_e^h denotes the spatial increment of size h in direction e , that is $D_e^h f(x) = f(x + he) - f(x)$. In [30] the authors study some properties of the set K of 2×2 matrices given by

$$K = \{(\Sigma_{e_1, e_2}(m), \Sigma_{\varepsilon_1, \varepsilon_2}(m)) : m \in S^1\} \subset \mathbb{R}^{2 \times 2}.$$

One of its key properties, obtained in [30, Lemma 7] and inspired from the work of Šverak on the Tartar conjecture [42], is the following inequality:

$$\det(X - Y) \gtrsim |X - Y|^4 \quad \forall (X, Y) \in K \times K.$$

Therefore the quantity Δ_{JK} can be estimated from below by

$$\Delta_{JK}(x, h, e) \gtrsim |D_e^h(\Sigma_{e_1, e_2}(m), \Sigma_{\varepsilon_1, \varepsilon_2}(m))(x)|^4 \gtrsim |D_e^h m(x)|^4,$$

where the last inequality follows from the (easily checkable) fact that $m \mapsto (\Sigma_{e_1, e_2}, \Sigma_{\varepsilon_1, \varepsilon_2})$ is an immersion.

Following [22], we aim to apply the div-curl Lemma, taking advantage of the fact that

$$\nabla \cdot \Sigma_{e_1, e_2}(m) = \mu_{e_1, e_2}, \quad \nabla \cdot \Sigma_{\varepsilon_1, \varepsilon_2}(m) = \mu_{\varepsilon_1, \varepsilon_2},$$

are locally finite measures. To this end let us fix χ a smooth cutoff function and set

$$E := \chi D_e^h \Sigma_{e_1, e_2}(m), \quad B := \chi D_e^h \Sigma_{\varepsilon_1, \varepsilon_2}(m).$$

Lemma 4.2. *For every $p \in]1, \infty[$ the following estimate holds true:*

$$\int_{\mathbb{R}^2} E \wedge B \, dx \lesssim pp' \|E\|_{L^p} \|\nabla \cdot B\|_{W^{-1, p'}} + \|B\|_{L^p} \|\nabla \cdot E\|_{W^{-1, p'}}. \quad (4.3)$$

Proof. The proof is nowadays standard, and we report it for the reader's convenience: for $1 < p < \infty$, using the potential theoretic solution ϕ to $\Delta\phi = \nabla \cdot E$, we find that E can be Hodge-decomposed as

$$E = \nabla\phi + \nabla^\perp\psi, \quad \|\nabla\phi\|_{L^{p'}(\mathbb{R}^2)} \lesssim pp' \|\nabla \cdot E\|_{W^{-1, p'}(\mathbb{R}^2)},$$

([24, Theorem 4.4.1], [43]), which yields

$$\begin{aligned} \left| \int_{\mathbb{R}^2} E \wedge B \, dx \right| &\leq \left| \int \nabla\phi \wedge B \right| + \left| \int \nabla^\perp\psi \wedge B \right| \\ &\lesssim \|\nabla\phi\|_{L^{p'}} \|B\|_{L^p} + \left| \int \left(\psi - \int_{\text{spt}(B)} \psi \right) \nabla \cdot B \right| \\ &\lesssim pp' \|\nabla \cdot E\|_{W^{-1, p'}(\mathbb{R}^2)} \|B\|_{L^p} + \|\nabla\psi\|_{L^p} \|\nabla \cdot B\|_{W^{-1, p'}(\mathbb{R}^2)} \\ &\lesssim pp' \|\nabla \cdot E\|_{W^{-1, p'}(\mathbb{R}^2)} \|B\|_{L^p} + (1 + pp') \|E\|_{L^p} \|\nabla \cdot B\|_{W^{-1, p'}(\mathbb{R}^2)}. \end{aligned}$$

The conclusion follows from $pp' \geq 4$. \square

Proposition 4.3. *Any solution m to (M) such that $\nabla \cdot \Sigma_{\alpha_1, \alpha_2}(m) \in \mathcal{M}_{\text{loc}}(\Omega)$ for $(\alpha_1, \alpha_2) = (e_1, e_2)$ and $(\varepsilon_1, \varepsilon_2)$, belongs to $B_{4, \infty; \text{loc}}^s(\Omega)$ for every $s < 4$.*

Proof. The div-curl estimate of Lemma 4.2 reduces the control of $\int \chi^2 |D_e^h m|^4$ to the estimate of the product

$$\|\chi D_e^h \Sigma_{e_1, e_2}(m)\|_{L^p} \|\nabla \cdot (\chi D_e^h \Sigma_{\varepsilon_1, \varepsilon_2}(m))\|_{W^{-1, p'}}$$

and its companion obtained by exchanging E and B . Let us for simplicity drop the frame index and write Σ instead of $\Sigma_{\alpha_1, \alpha_1}$, and also write D^h instead of D_e^h . We start by estimating the $W^{-1, p'}$ norm of the second factor

$$\nabla \cdot (\chi D_e^h \Sigma(m)) = \chi D^h \mu + D_e^h \Sigma(m) \cdot \nabla \chi.$$

By Sobolev embedding $W^{1, p} \subset C^{1 - \frac{2}{p}}$ for $p > 2$, it holds

$$\begin{aligned} \|\chi D^h \mu\|_{W^{-1, p'}} &= \sup \left\{ \int D^h(\chi\psi) d\mu, \|\psi\|_{W^{1, p}} \leq 1 \right\} \\ &\leq \sup_{\|\psi\|_{W^{1, p}} \leq 1} \|\mu\|_{\mathcal{M}} \|D^h(\chi\psi)\|_{\infty} \\ &\lesssim \|\mu\|_{\mathcal{M}} |h|^{1 - \frac{2}{p}}, \end{aligned}$$

and therefore

$$\|\nabla \cdot (\chi D_e^h \Sigma(m))\|_{W^{-1, p'}} \lesssim \|\mu\|_{\mathcal{M}} |h|^{1 - \frac{2}{p}} + |h|.$$

The L^p norms $\|\chi D_e^h \Sigma\|_{L^p}$ are uniformly bounded. Inserting this estimate in (4.3), and choosing $p = -\log(|h|)$ [43], one easily obtains the modulus of continuity

$$\int \chi^2 |D_e^h m|^4 \, dx \lesssim \int E \wedge B \, dx \lesssim (1 + \|\mu\|_{\mathcal{M}}) |h| \log\left(\frac{1}{|h|}\right)$$

for $|h| < \exp(-2)$. This implies $m \in B_{4, \infty}^s(\text{spt}\chi)$ for every $s < \frac{1}{4}$. \square

Remark 4.4. It is unclear to the authors whether the $\frac{1}{4}$ exponent is optimal or not.

BIBLIOGRAPHY

- [1] ALOUGES, F., RIVIÈRE, T., AND SERFATY, S. Néel and cross-tie wall energies for planar micromagnetic configurations. *ESAIM Control Optim. Calc. Var.* 8 (2002), 31–68. A tribute to J. L. Lions.
- [2] AMBROSIO, L., DE LELLIS, C., AND MANTEGAZZA, C. Line energies for gradient vector fields in the plane. *Calc. Var. Partial Differential Equations* 9, 4 (1999), 327–255.
- [3] AMBROSIO, L., FUSCO, N., AND PALLARA, D. *Functions of bounded variation and free discontinuity problems*. Oxford Mathematical Monographs. The Clarendon Press, Oxford University Press, New York, 2000.
- [4] AMBROSIO, L., KIRCHHEIM, B., LECUMBERRY, M., AND RIVIÈRE, T. On the rectifiability of defect measures arising in a micromagnetics model. In *Nonlinear problems in mathematical physics and related topics, II*, vol. 2 of *Int. Math. Ser. (N. Y.)*. Kluwer/Plenum, New York, 2002, pp. 29–60.
- [5] AVILES, P., AND GIGA, Y. A mathematical problem related to the physical theory of liquid crystal configurations. In *Miniconference on geometry and partial differential equations, 2 (Canberra, 1986)*, vol. 12 of *Proc. Centre Math. Anal. Austral. Nat. Univ.* Austral. Nat. Univ., Canberra, 1987, pp. 1–16.
- [6] AVILES, P., AND GIGA, Y. On lower semicontinuity of a defect energy obtained by a singular limit of the Ginzburg-Landau type energy for gradient fields. *Proc. Roy. Soc. Edinburgh Sect. A* 129, 1 (1999), 1–17.
- [7] BAHOURI, H., CHEMIN, J.-Y., AND DANCHIN, R. *Fourier analysis and nonlinear partial differential equations*, vol. 343 of *Grundlehren der Mathematischen Wissenschaften [Fundamental Principles of Mathematical Sciences]*. Springer, Heidelberg, 2011.
- [8] BELLETTINI, G., BERTINI, L., MARIANI, M., AND NOVAGA, M. Γ -entropy cost for scalar conservation laws. *Arch. Ration. Mech. Anal.* 195, 1 (2010), 261–309.
- [9] BOCHARD, P., AND IGNAT, R. Kinetic formulation of vortex vector fields. *Anal. PDE* 10, 3 (2017), 729–756.
- [10] CHESKIDOV, A., CONSTANTIN, P., FRIEDLANDER, S., AND SHVYDKOY, R. Energy conservation and Onsager’s conjecture for the Euler equations. *Nonlinearity* 21, 6 (2008), 1233–1252.
- [11] CONSTANTIN, P., E, W., AND TITI, E. S. Onsager’s conjecture on the energy conservation for solutions of Euler’s equation. *Comm. Math. Phys.* 165, 1 (1994), 207–209.
- [12] CONTI, S., AND DE LELLIS, C. Sharp upper bounds for a variational problem with singular perturbation. *Math. Ann.* 338, 1 (2007), 119–146.
- [13] CRASTA, G., DE CICCO, V., DE PHILIPPIS, G., AND GHIRALDIN, F. Structure of solutions of multidimensional conservation laws with discontinuous flux and applications to uniqueness. *Arch. Ration. Mech. Anal.* 221, 2 (2016), 961–985.
- [14] CRIPPA, G., OTTO, F., AND WESTDICKENBERG, M. Regularizing effect of nonlinearity in multidimensional scalar conservation laws. In *Transport equations and multi-D hyperbolic conservation laws*, vol. 5 of *Lect. Notes Unione Mat. Ital.* Springer, Berlin, 2008, pp. 77–128.
- [15] DE LELLIS, C., AND IGNAT, R. A regularizing property of the 2D-eikonal equation. *Comm. Partial Differential Equations* 40, 8 (2015), 1543–1557.
- [16] DE LELLIS, C., AND OTTO, F. Structure of entropy solutions to the eikonal equation. *J. Eur. Math. Soc. (JEMS)* 5, 2 (2003), 107–145.
- [17] DE LELLIS, C., OTTO, F., AND WESTDICKENBERG, M. Structure of entropy solutions for multi-dimensional scalar conservation laws. *Arch. Ration. Mech. Anal.* 170, 2 (2003), 137–184.
- [18] DE LELLIS, C., AND SZÉKELYHIDI, L. On h -principle and Onsager’s conjecture. *Eur. Math. Soc. Newsl.*, 95 (2015), 19–24.
- [19] DE LELLIS, C., AND WESTDICKENBERG, M. On the optimality of velocity averaging lemmas. *Ann. Inst. H. Poincaré Anal. Non Linéaire* 20, 6 (2003), 1075–1085.
- [20] DESIMONE, A., MÜLLER, S., KOHN, R., AND OTTO, F. A compactness result in the gradient theory of phase transitions. *Proc. Roy. Soc. Edinburgh Sect. A* 131, 4 (2001), 833–844.
- [21] GOLDMAN, M., JOSIEN, M., AND OTTO, F. New bounds for the inhomogenous Burgers and the Kuramoto-Sivashinsky equations. *Comm. Partial Differential Equations* 40, 12 (2015), 2237–2265.
- [22] GOLSE, F. Nonlinear regularizing effect for hyperbolic partial differential equations. In *XVIth International Congress on Mathematical Physics*. World Sci. Publ., Hackensack, NJ, 2010, pp. 433–437.
- [23] GOLSE, F., AND PERTHAME, B. Optimal regularizing effect for scalar conservation laws. *Rev. Mat. Iberoam.* 29, 4 (2013), 1477–1504.

- [24] GRAFAKOS, L. *Classical Fourier analysis*, second ed., vol. 249 of *Graduate Texts in Mathematics*. Springer, New York, 2008.
- [25] IGNAT, R., AND MERLET, B. Entropy method for line-energies. *Calc. Var. Partial Differential Equations* 44, 3-4 (2012), 375–418.
- [26] JABIN, P., OTTO, F., AND PERTHAME, B. Line-energy Ginzburg-Landau models: zero-energy states. *Ann. Sc. Norm. Super. Pisa Cl. Sci. (5) 1*, 1 (2002), 187–202.
- [27] JABIN, P.-E., AND PERTHAME, B. Compactness in Ginzburg-Landau energy by kinetic averaging. *Comm. Pure Appl. Math.* 54, 9 (2001), 1096–1109.
- [28] JABIN, P.-E., AND PERTHAME, B. Regularity in kinetic formulations via averaging lemmas. *ESAIM Control Optim. Calc. Var.* 8 (2002), 761–774. A tribute to J. L. Lions.
- [29] JIN, W., AND KOHN, R. V. Singular perturbation and the energy of folds. *J. Nonlinear Sci.* 10, 3 (2000), 355–390.
- [30] LORENT, A., AND PENG, G. Regularity of the Eikonal equation with two vanishing entropies. *Ann. Inst. H. Poincaré Anal. Non Linéaire* (2017).
- [31] ORTIZ, M., AND GIOIA, G. The morphology and folding patterns of buckling-driven thin-film blisters. *J. Mech. Phys. Solids* 42, 3 (1994), 531–559.
- [32] OTTO, F., AND STEINER, J. The concertina pattern: from micromagnetics to domain theory. *Calc. Var. Partial Differential Equations* 39, 1-2 (2010), 139–181.
- [33] PERTHAME, B. *Kinetic formulation of conservation laws*, vol. 21 of *Oxford Lecture Series in Mathematics and its Applications*. Oxford University Press, Oxford, 2002.
- [34] PERTHAME, B. T. *Kinetic formulation of conservation laws*, vol. 21 of *Oxford Lecture Series in Mathematics and its Applications*. Oxford University Press, Oxford, 2002.
- [35] POLIAKOVSKY, A. Jumps detection in Besov spaces via a new BBM formula. applications to Aviles-Giga type functionals. arXiv:1703.04208, 2017.
- [36] RIVIÈRE, T., AND SERFATY, S. Limiting domain wall energy for a problem related to micromagnetics. *Comm. Pure Appl. Math.* 54, 3 (2001), 294–338.
- [37] RIVIÈRE, T., AND SERFATY, S. Compactness, kinetic formulation, and entropies for a problem related to micromagnetics. *Comm. Partial Differential Equations* 28, 1-2 (2003), 249–269.
- [38] TARTAR, L. Compensated compactness and applications to partial differential equations. In *Nonlinear analysis and mechanics: Heriot-Watt Symposium, Vol. IV*, vol. 39 of *Res. Notes in Math*. Pitman, Boston, Mass.-London, 1979, pp. 136–212.
- [39] TARTAR, L. *From hyperbolic systems to kinetic theory*, vol. 6 of *Lecture Notes of the Unione Matematica Italiana*. Springer-Verlag, Berlin; UMI, Bologna, 2008. A personalized quest.
- [40] TRIEBEL, H. *Theory of function spaces. III*, vol. 100 of *Monographs in Mathematics*. Birkhäuser Verlag, Basel, 2006.
- [41] VASSEUR, A. Strong traces for solutions of multidimensional scalar conservation laws. *Arch. Ration. Mech. Anal.* 160, 3 (2001), 181–193.
- [42] ŠVERÁK, V. On Tartar’s conjecture. *Ann. Inst. H. Poincaré Anal. Non Linéaire* 10, 4 (1993), 405–412.
- [43] YUDOVICH, V. I. Uniqueness theorem for the basic nonstationary problem in the dynamics of an ideal incompressible fluid. *Math. Res. Lett.* 2, 1 (1995), 27–38.

F.G.: UNIVERSITÄT BASEL, SPIEGELGASSE 1, CH-4051 BASEL, SWITZERLAND

E-mail address: Francesco.Ghiraldin@unibas.ch

X.L.: INSTITUT DE MATHÉMATIQUES DE TOULOUSE, CNRS UMR 5219, UNIVERSITÉ PAUL SABATIER, TOULOUSE, FRANCE

E-mail address: Xavier.Lamy@math.univ-toulouse.fr

LATEST PREPRINTS

- | No. | Author: Title |
|---------|--|
| 2016-31 | M. Dambrine, H. Harbrecht, M. D. Peters, B. Puig
<i>On Bernoulli's free boundary problem with a random boundary</i> |
| 2016-32 | H. Harbrecht, J. Tausch
<i>A fast sparse grid based space-time boundary element method for the nonstationary heat equation</i> |
| 2016-33 | S. Iula
<i>A note on the Moser-Trudinger inequality in Sobolev-Slobodeckij spaces in dimension one</i> |
| 2016-34 | C. Bürli, H. Harbrecht, P. Odermatt, S. Sayasone, N. Chitnis
<i>Mathematical analysis of the transmission dynamics of the liver fluke, <i>Opisthorchis viverrini</i></i> |
| 2017-01 | J. Dölz and T. Gerig, M. Lüthi, H. Harbrecht and T. Vetter
<i>Efficient computation of low-rank Gaussian process models for surface and image registration</i> |
| 2017-02 | M. J. Grote, M. Mehlin, S. A. Sauter
<i>Convergence analysis of energy conserving explicit local time-stepping methods for the wave equation</i> |
| 2017-03 | Y. Bilu, F. Luca, D. Masser
<i>Collinear CM-points</i> |
| 2017-04 | P. Zaspel
<i>Ensemble Kalman filters for reliability estimation in perfusion inference</i> |
| 2017-05 | J. Dölz and H. Harbrecht
<i>Hierarchical Matrix Approximation for the Uncertainty Quantification of Potentials on Random Domains</i> |
| 2017-06 | P. Zaspel
<i>Analysis and parallelization strategies for Ruge-Stüben AMG on many-core processor</i> |
| 2017-07 | H. Harbrecht and M. Schmidlin
<i>Multilevel Methods for Uncertainty Quantification of Elliptic PDEs with Random Anisotropic Diffusion</i> |
| 2017-08 | M. Griebel and H. Harbrecht
<i>Singular value decomposition versus sparse grids: Refined complexity Estimates</i> |

LATEST PREPRINTS

- No.** **Author:** *Title*
- 2017-09 **J. Garcke and I. Kalmykov**
Efficient Higher Order Time Discretization Schemes for Hamilton-Jacobi-Bellman Equations Based on Diagonally Implicit Symplectic Runge-Kutta Methods
- 2017-10 **M. J. Grote and U. Nahum**
Adaptive Eigenspace Regularization For Inverse Scattering Problems
- 2017-11 **J. Dölz, H. Harbrecht, S. Kurz, S. Schöps and F. Wolf**
A Fast Isogeometric BEM for the Three Dimensional Laplace- and Helmholtz Problems
- 2017-12 **P. Zaspel**
Algorithmic patterns for H -matrices on many-core processors
- 2017-13 **R. Brügger, R. Croce and H. Harbrecht**
Solving a free boundary problem with non-constant coefficients
- 2017-14 **M. Dambrine, H. Harbrecht and B. Puig**
Incorporating knowledge on the measurement noise in electrical impedance tomography
- 2017-15 **C. Bürli, H. Harbrecht, P. Odermatt, S. Sayasone and N. Chitnis**
*Analysis of Interventions against the Liver Fluke, *Opisthorchis viverrini**
- 2017-16 **D. W. Masser**
Abcological anecdotes
- 2017-17 **P. Corvaja, D. W. Masser and U. Zannier**
Torsion hypersurfaces on abelian schemes and Betti coordinates
- 2017-18 **F. Caubet, M. Dambrine and H. Harbrecht**
A Newton method for the data completion problem and application to obstacle detection in Electrical Impedance Tomography
- 2018-01 **H. Harbrecht and P. Zaspel**
On the algebraic construction of sparse multilevel approximations of elliptic tensor product problems
- 2018-02 **F. Ghiraldin and X. Lamy**
Optimal Besov differentiability for entropy solutions of the eikonal equation